

# Advanced QFT: Problem Set 2

February 19, 2019

## Problem 1 *Dirac equation*

- a) Demonstrate the Lorentz invariance of the Dirac equation.
- b) Check whether  $\psi^\dagger(x)\psi(x)$  and  $\bar{\psi}(x)\psi(x)$  are Lorentz invariant.
- c) Show that  $\sum_r u_r(p)\bar{u}_r(p) = \not{p} + m$  and that  $\sum_r v_r(p)\bar{v}_r(p) = \not{p} - m$ .

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## Problem 2 *Grassman variables*

Evaluate the following Grassman integrals

- a)  $\int d^N \theta^* d^N \theta \times e^{-\bar{\theta}^\dagger \bar{\theta}} = 1$
- b)  $\frac{1}{\det(A)} \int d^N \theta'^* d^N \theta' \times e^{-\bar{\theta}'^\dagger A \bar{\theta}'} = 1$
- c)  $\int d^N \theta^* d^N \theta \times e^{-\bar{\theta}^\dagger A \bar{\theta} + \bar{K}^\dagger \bar{\theta} + \bar{\theta}^\dagger \bar{K}} = \det(A) e^{\bar{K}^\dagger A^{-1} \bar{K}}$ . (1)
- d)  $\int d^N \theta^* d^N \theta \times \theta_k \theta_l^* e^{-\bar{\theta}^\dagger A \bar{\theta}} = \det(A) (A^{-1})_{kl}$ .

where  $\{\theta_i\}$  and  $\{\theta_i^*\}$  are two independent sets of Grassmann numbers and A is a c-number skew-symmetric  $N \times N$  matrix. Notice that the determinant is in the numerator, whereas for c-number Gaussian integrals

$$\int_{-\infty}^{-\infty} \exp\left(-\frac{1}{2} x^T A x\right) d^n x = \sqrt{\frac{(2\pi)^n}{\det(A)}}$$

List of conventions

$$\begin{aligned} \int d\theta \theta &= 1 \\ f(\theta + d\theta) &= f(\theta) + d\theta f'(\theta) \quad (\text{left derivative}) \\ d^N \theta &= d\theta_N d\theta_{N-1} \dots d\theta_1 \quad (\text{descending order}) \\ d^N \theta^* d^N \theta &= d\theta_N^* d\theta_N \dots d\theta_1^* d\theta_1 \quad (\text{descending order}) \\ f(\theta) &= f_0 + \theta_i f_1^i + \dots + \theta_1 \theta_2 \dots \theta_N f_N; \quad (\text{ascending order}) \end{aligned}$$

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**Problem 3** *Gamma matrix manipulation*

Consider the matrix  $\gamma^5 = i\gamma^0\gamma^1\gamma^2\gamma^3$ .

- a) Show that  $\gamma^5$  anticommutes with each of the  $\gamma^\mu$  matrices,  $\gamma^5\gamma^\mu = -\gamma^\mu\gamma^5$ .
- b) Show that  $\gamma^5$  is hermitian and that  $(\gamma^5)^2 = 1$ .
- c) Show that  $\gamma^5 = (-i/24)\epsilon_{\kappa\lambda\mu\nu}\gamma^\kappa\gamma^\lambda\gamma^\mu\gamma^\nu$ .  
(Sign convention:  $\epsilon^{0123} = +1$ ,  $\epsilon_{0123} = -1$ .)
- d) Show that  $\gamma^{[\lambda}\gamma^\mu\gamma^{\nu]} = -i\epsilon^{\kappa\lambda\mu\nu}\gamma_\kappa\gamma^5$ .
- e) Show that any  $4 \times 4$  matrix  $\Gamma$  is a unique linear combination of the following 16 matrices:  $1$ ,  $\gamma^\mu$ ,  $\gamma^{[\mu}\gamma^{\nu]}$ ,  $\gamma^5\gamma^\mu$  and  $\gamma^5$ .
- f) Using  $\{\gamma^\mu, \gamma^\nu\} = 2g^{\mu\nu}$  show that in  $d$  dimensions

$$\begin{aligned} \gamma^\mu\gamma_\mu &= dI_4, \\ \gamma^\mu\gamma^\nu\gamma_\mu &= -(d-2)\gamma^\nu, \\ \gamma^\mu\gamma^\nu\gamma^\rho\gamma_\mu &= 2\{\gamma^\rho, \gamma^\nu\} + (d-4)\gamma^\nu\gamma^\rho = 4g^{\nu\rho}I_4 - (4-d)\gamma^\nu\gamma^\rho, \\ \gamma^\mu\gamma^\nu\gamma^\rho\gamma^\sigma\gamma_\mu &= -2\gamma^\sigma\gamma^\rho\gamma^\nu + (4-d)\gamma^\nu\gamma^\rho\gamma^\sigma, \\ \not{p}^2 &= p^2. \end{aligned}$$

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**Problem 4** *Anomalous magnetic moment of the electron*

a) Draw the Feynman diagram that represents the leading order contribution to the interaction between a photon and an electron. Write down the amplitude for this diagram.

b) Evaluating  $\bar{u}(p', s') (\not{p}'\gamma^\mu + \gamma^\mu\not{p}) u(p, s)$  in two different ways, prove the Gordon identity

$$\bar{u}(p', s')\gamma^\mu u(p, s) = \frac{1}{2m}\bar{u}(p', s')\left[(p+p')^\mu + i\sigma^{\mu\nu}(p'-p)_\nu\right]u(p, s).$$

c) Prove that the most general matrix element of the electromagnetic current allowed by Lorentz invariance and current conservation is

$$\begin{aligned} \langle p', s' | J^\mu(x) | p, s \rangle &= \bar{u}(p', s') \left[ \gamma^\mu F_1(q^2) + \frac{i\sigma^{\mu\nu}(p'-p)_\nu}{2m} F_2(q^2) \right] u(p, s) \\ &= \bar{u}(p', s') \left[ \frac{(p'+p)^\mu}{2m} F_1(q^2) + \frac{i\sigma^{\mu\nu}(p'-p)_\nu}{2m} (F_1(q^2) + F_2(q^2)) \right] u(p, s), \end{aligned} \tag{-2}$$

where  $F_1(q^2)$  and  $F_2(q^2)$  are two arbitrary functions of the momentum transfer,  $q^\mu = p^\mu - p'^\mu$ , called form factors.  $|p, s\rangle$  is a one-particle state of an electron with momentum  $p^\mu$  and spin  $s$ . Notice that  $\langle p', s' | J^\mu(x) | p, s \rangle$  transforms as a vector under Lorentz transformations, so a good start is to write down all possible vector interactions. Using translation invariance

$$\langle p', s' | J^\mu(x) | p, s \rangle = \langle p', s' | J^\mu(0) | p, s \rangle e^{i(p'-p)x},$$

and the conservation of the electromagnetic current,  $\partial_\mu J^\mu = 0$ , show that

$$(p' - p)_\mu \langle p', s' | J^\mu(0) | p, s \rangle = 0$$

The coefficient of the first term in eq(??) is what we observe as an electric charge. Therefore it is natural to absorb  $F_1(d^2)$  into the definition of the electric charge  $e$ , i.e. we set  $F_1(0) = 1$  in the on-shell scheme. On the other hand, the second term is related to the magnetic moment of the electron, which gets shifted from the tree-level value by  $1 + F_2(0)$ , which is now to be calculated to one-loop order.

**d)** Draw the Feynman diagrams that contribute to next-to-leading order to the interaction between a photon and an electron. Explain why only one of these diagrams contributes to  $F_2(q^2)$  and evaluate the diagram. In the calculation, concentrate on those terms that contribute to  $F_2(q^2)$ , which are finite and do not need any regulator, and drop the other contributions. Evaluate your result for  $F_2(0)$  numerically, and compare it to the experimental value for  $(g_e - 2)$ .

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