

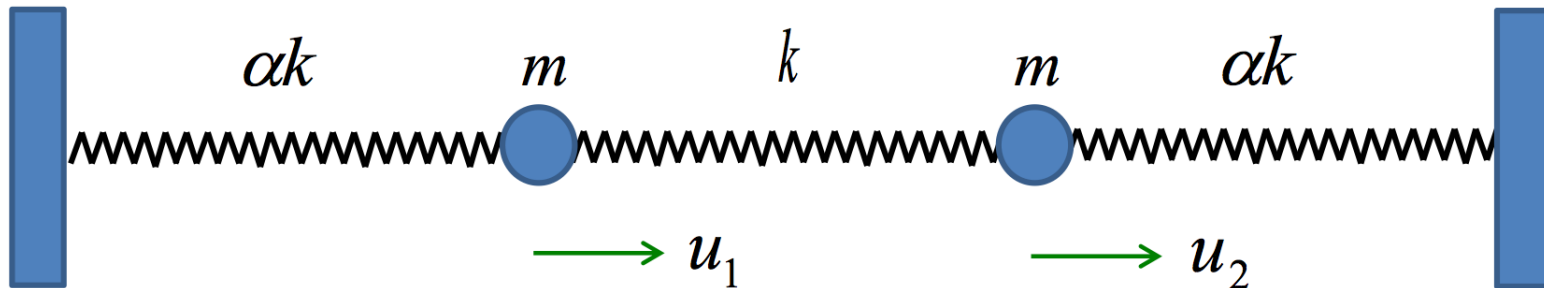
# Spring-mass systems

Now consider a horizontal system in the form of masses on springs

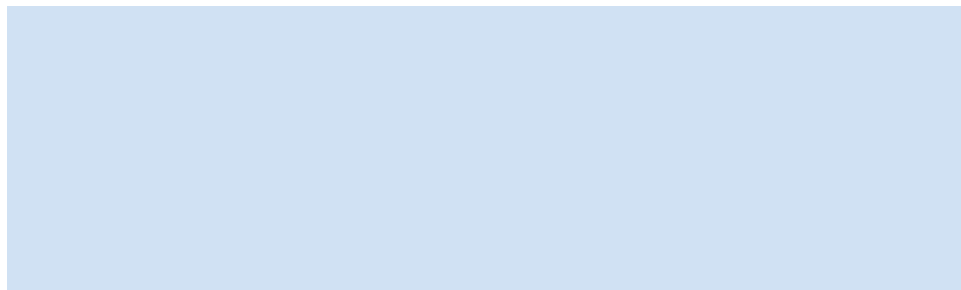
- Again solve via decoupling and matrix methods
- Obtain the energy within the system
- Find specific solutions

# Horizontal spring-mass system

Consider two masses moving without friction, between three springs, two with spring constants  $\alpha k$ , one with spring constant  $k$

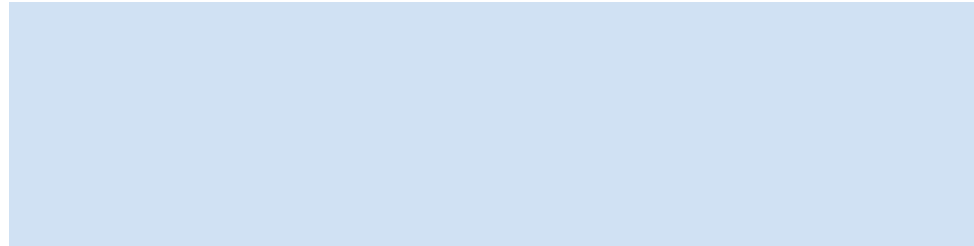


Equations of motion:



# Solutions of horizontal spring-mass system

Equations of motion:



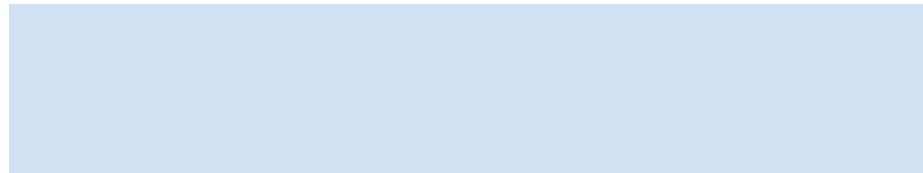
(1)

(2)

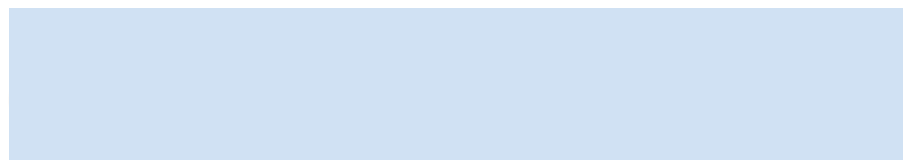
Solve by decoupling method (add 1 and 2 and subtract 2 from 1). As before, we can write down the normal coordinates, call them  $q_1$  and  $q_2$

$$q_1 = \frac{1}{\sqrt{2}}(u_1 + u_2) \quad \text{which means...} \quad u_1 = \frac{1}{\sqrt{2}}(q_1 + q_2)$$
$$q_2 = \frac{1}{\sqrt{2}}(u_1 - u_2) \quad u_2 = \frac{1}{\sqrt{2}}(q_1 - q_2)$$

Substituting gives:



Gives normal frequencies of:



Centre of Mass

Relative

# Cross-checking with matrix method

Write equations of motion as homogenous matrix equation

$$\begin{pmatrix} \frac{d^2}{dt^2} + \frac{\alpha k}{m} + \frac{k}{m} & -\frac{k}{m} \\ -\frac{k}{m} & \frac{d^2}{dt^2} + \frac{\alpha k}{m} + \frac{k}{m} \end{pmatrix} \begin{pmatrix} u_1 \\ u_2 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

Demand that the resulting operator matrix is singular, i.e.  $\text{Det}\{\text{matrix}\}=0$

$$\begin{vmatrix} -\omega^2 + \frac{\alpha k}{m} + \frac{k}{m} & -\frac{k}{m} \\ -\frac{k}{m} & -\omega^2 + \frac{\alpha k}{m} + \frac{k}{m} \end{vmatrix} = 0$$

We obtain same solutions as before

Substitute in the below trial solution

$$\begin{pmatrix} u_1 \\ u_2 \end{pmatrix} = \text{Re} \begin{pmatrix} X \\ Y \end{pmatrix} e^{i\omega t}$$

Hence get eigenvalue equation

$$\left( -\omega^2 + \frac{\alpha k}{m} + \frac{k}{m} \right)^2 - \left( \frac{k}{m} \right)^2 = 0$$

# Energy of horizontal spring-mass system

Kinetic Energy

$$K = \frac{1}{2}m(\dot{u}_1^2 + \dot{u}_2^2)$$

$$K = \frac{1}{2}m(\dot{q}_1^2 + \dot{q}_2^2)$$

Potential Energy

$$V = \frac{1}{2}\alpha k u_1^2 + \frac{1}{2}k(u_2 - u_1)^2 + \frac{1}{2}\alpha k u_2^2$$

$$V = \frac{1}{2}\alpha k q_1^2 + \frac{1}{2}(\alpha + 2)k q_2^2$$

$$V = \frac{1}{2}m\omega_1^2 q_1^2 + \frac{1}{2}m\omega_2^2 q_2^2$$

Total Energy:



Sum of energies in each normal mode

# Specific Solutions for horizontal spring-mass system

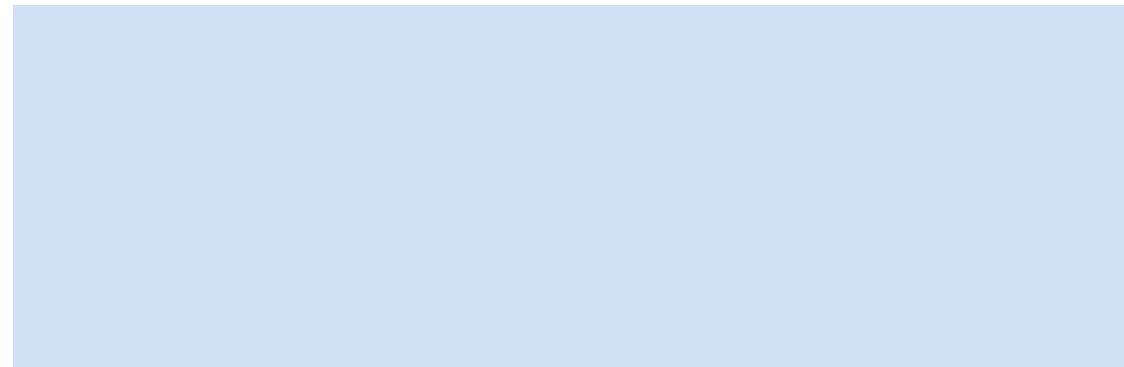
General solution is the sum of the two normal modes

$$u_1 = A_1 \cos(\omega_1 t + \phi_1) + A_2 \cos(\omega_2 t + \phi_2)$$

$$u_2 = A_1 \cos(\omega_1 t + \phi_1) - A_2 \cos(\omega_2 t + \phi_2)$$

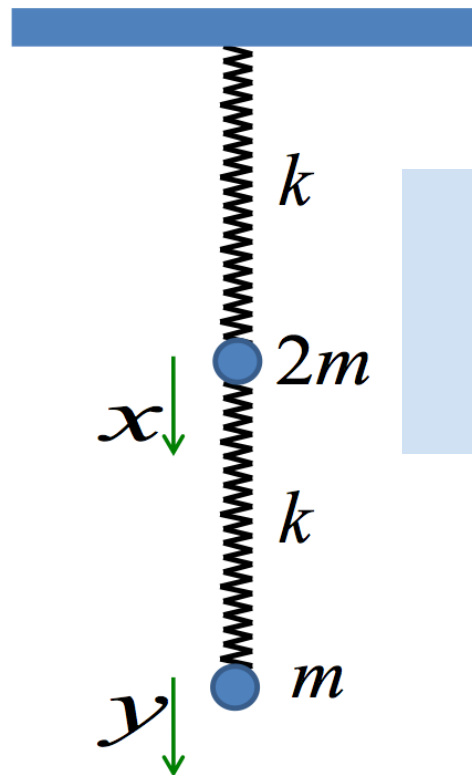
Initial conditions:  $u_1 = a$  ;  $u_2 = 0$  ;  $\dot{u}_1 = \dot{u}_2 = 0$

These give:  $A_1 = A_2 = \frac{a}{2}$  ;  $\phi_1 = \phi_2 = 0$



Both normal modes excited. The 'beats' solution.  
Completely analagous to the coupled pendulum

# Vertical spring-mass system



$x$  and  $y$  are displacements  
from equilibrium positions

- Find the normal frequencies of the system
- Find the ratio of the amplitudes for each normal mode

Decoupling method only works for limited cases with a sufficient amount of symmetry. You cannot solve this with decoupling, so have to go to matrix method and have a guess!

# Solving with matrix method

Write equations of motion as homogenous matrix equation

$$\begin{pmatrix} \frac{d^2}{dt^2} + \frac{k}{m} & -\frac{k}{2m} \\ -\frac{k}{m} & \frac{d^2}{dt^2} + \frac{k}{m} \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

Demand that the resulting operator matrix is singular, i.e.  $\text{Det}\{\text{matrix}\}=0$

$$\begin{vmatrix} -\omega^2 + \frac{k}{m} & -\frac{k}{2m} \\ -\frac{k}{m} & -\omega^2 + \frac{k}{m} \end{vmatrix} = 0$$

From this we obtain the normal frequencies

Substitute in the below trial solution

$$\begin{pmatrix} x \\ y \end{pmatrix} = \text{Re} \begin{pmatrix} X \\ Y \end{pmatrix} e^{i\omega t}$$

Hence get eigenvalue equation

$$\left( -\omega^2 + \frac{k}{m} \right)^2 - \frac{1}{2} \left( \frac{k}{m} \right)^2 = 0$$



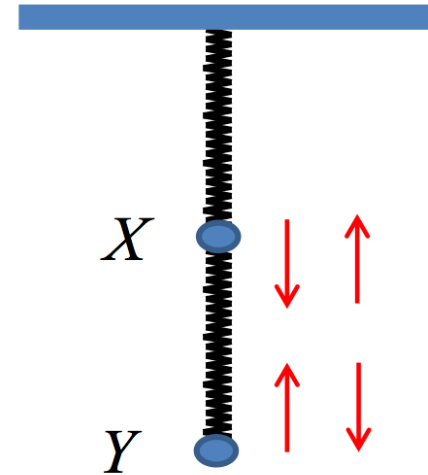
# Normal modes of vertical spring-mass system

Normal mode 1

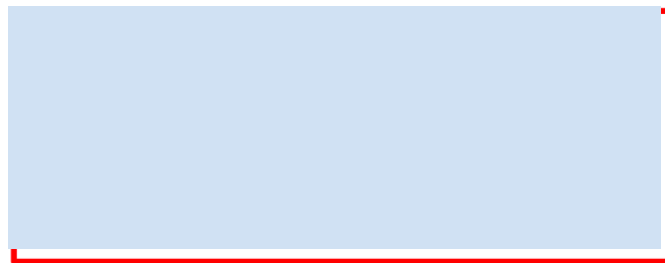


Substitute back into eigenvector equation  $\left(-\omega_{1,2}^2 + \frac{k}{m}\right)X - \left(\frac{k}{2m}\right)Y = 0$

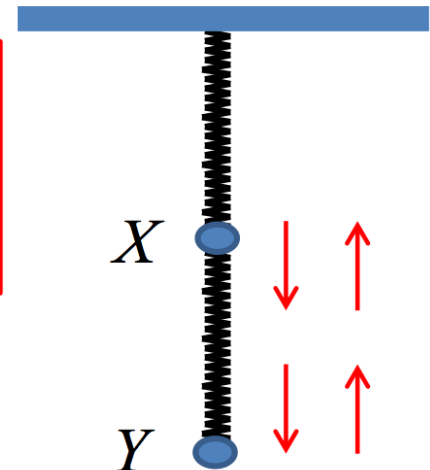
to yield  $X/Y = -1/\sqrt{2}$



Normal mode 2



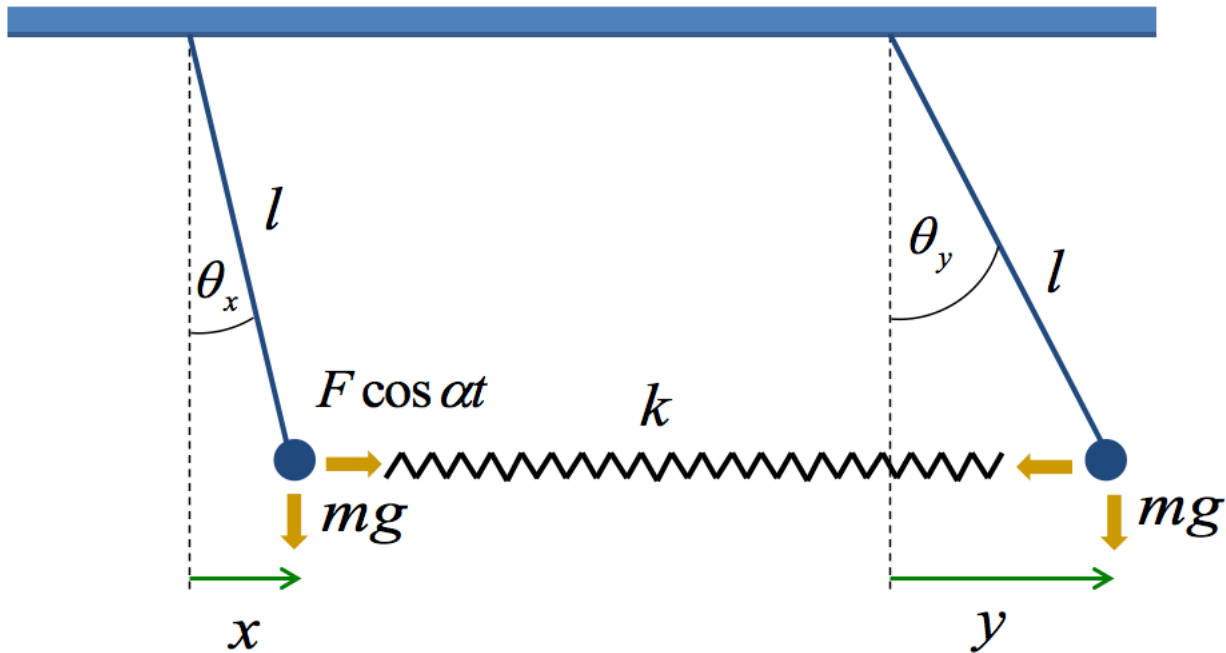
corresponds to  $X/Y = 1/\sqrt{2}$



# Coupled Oscillators with a driving force

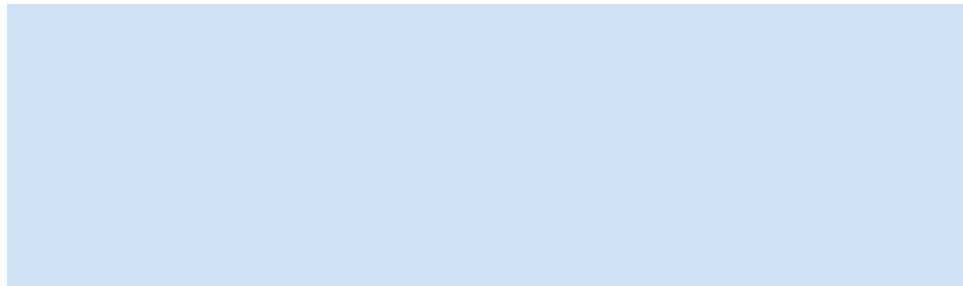
- So the last physical system we are going to look at in this first part of the course is the forced coupled pendula, along with a damping factor
  1. Finding the Complementary Function
  2. Finding the particular integral
- Then do the same for a horizontal spring-mass system

# Damped driven coupled pendula



Both pendula experience a retarding force of  $\gamma \times$  velocity

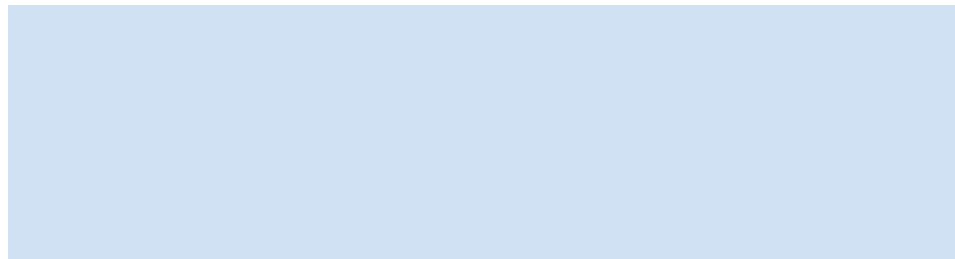
Equations of motion:



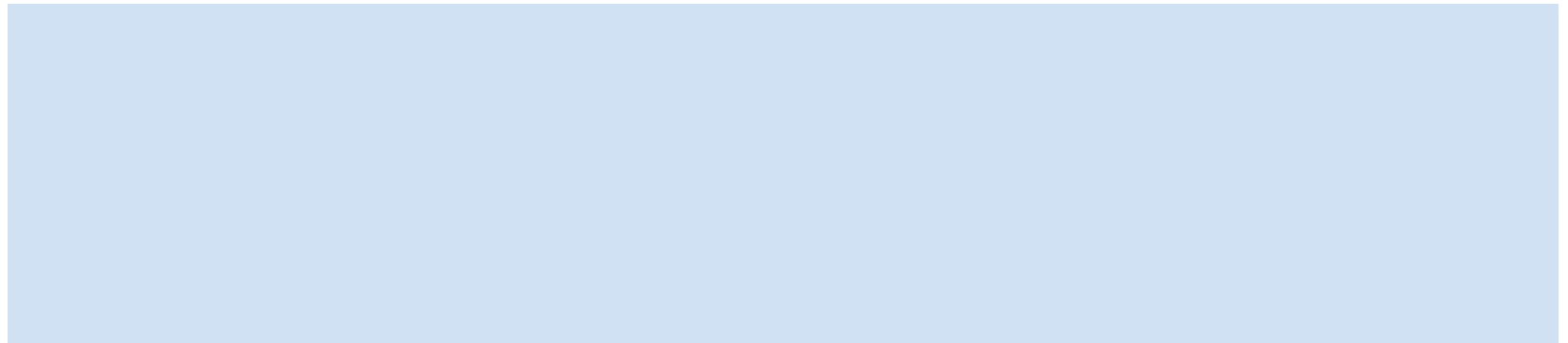
# Damped driven coupled pendula

Let's arrange equations of motion in form  $A \begin{pmatrix} x \\ y \end{pmatrix} = \text{whatever}$

We have:



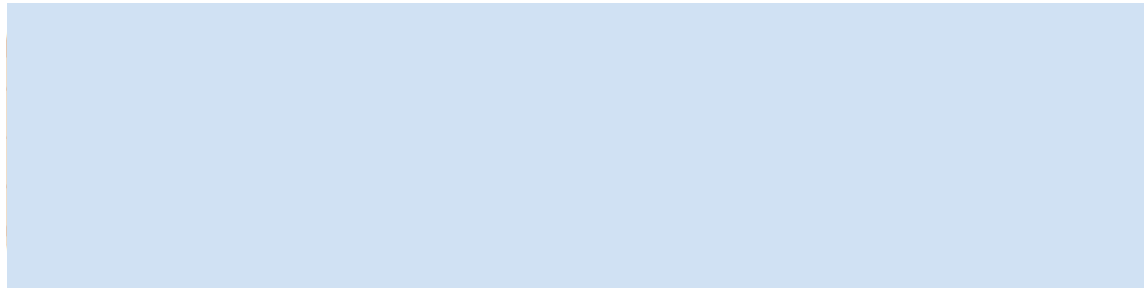
and so



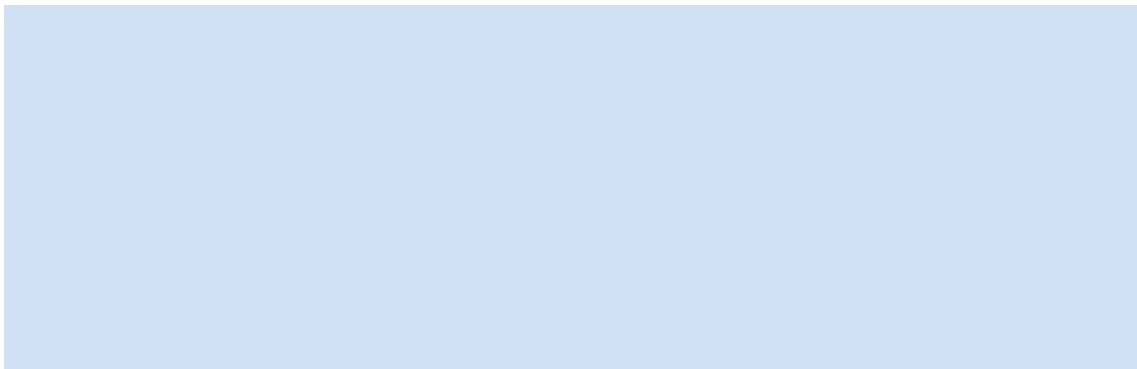
Contrary to before, this equation is inhomogeneous, in that  $\text{RHS} \neq 0$ .  
To solve it we need to find both the complementary function (CF), which is solution to the homogeneous equivalent, and the particular integral (PI)

# Damped driven coupled pendula: finding CF

To find CF write down homogenous equation and solve as previously

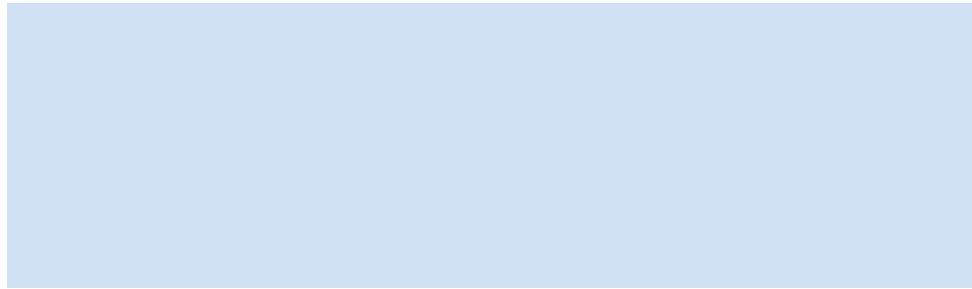


Try  $\begin{pmatrix} x \\ y \end{pmatrix} = \text{Re} \left( \begin{pmatrix} X \\ Y \end{pmatrix} e^{i\omega t} \right)$  and find solution when operator matrix is singular


$$= 0$$

# Damped driven coupled pendula: finding CF

To find CF write down homogenous equation and solve as previously


$$= 0$$

$$\Rightarrow \left[ -\omega^2 + \frac{i\gamma}{m}\omega + \left( \frac{g}{l} + \frac{k}{m} \right) \right]^2 - \left( \frac{k}{m} \right)^2 = 0 \quad \Rightarrow \left[ -\omega^2 + \frac{i\gamma}{m}\omega + \left( \frac{g}{l} + \frac{k}{m} \right) \right] = \pm \left( \frac{k}{m} \right)$$

gives the equations

$$\bar{\omega}_1^2 - \frac{i\gamma}{m}\omega - \frac{g}{l} = 0 \quad \text{and} \quad \bar{\omega}_2^2 - \frac{i\gamma}{m}\omega - \left( \frac{g}{l} + 2\frac{k}{m} \right) = 0$$

with solutions



where

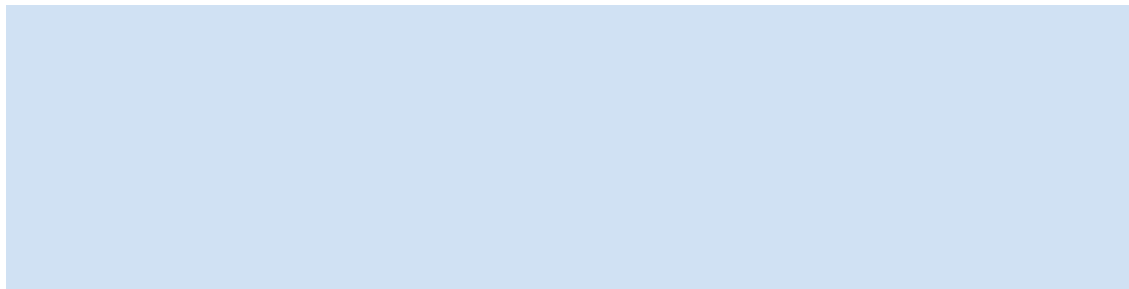
$$\omega_1^2 = \frac{g}{l}$$
$$\omega_2^2 = \frac{g}{l} + 2\frac{k}{m}$$

which are the results from the undamped scenario

There is no physical difference between +/- variants. Just use + from now on.

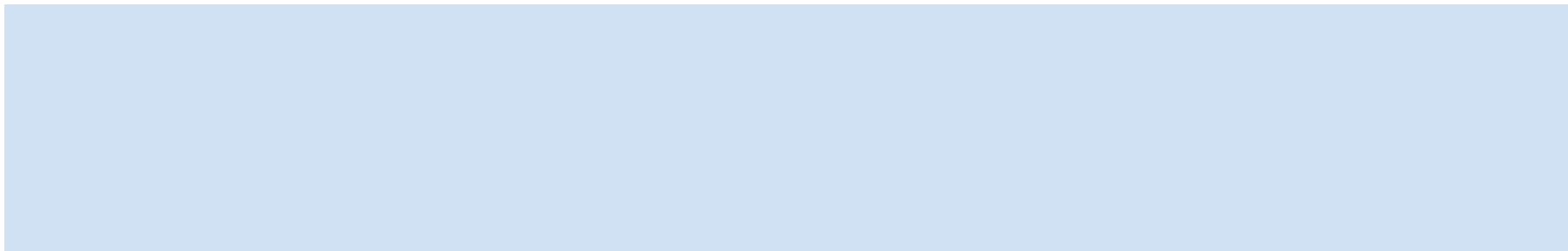
# Damped driven coupled pendula: finding CF

Substitute eigenvalues into the below:


$$\begin{pmatrix} X \\ Y \end{pmatrix} = 0$$

to deduce one mode has  $X=Y$ , & the other  $X=-Y$ . Since  $\begin{pmatrix} x \\ y \end{pmatrix} = \text{Re} \begin{pmatrix} X \\ Y \end{pmatrix} \exp(i\omega_{1,2}t)$

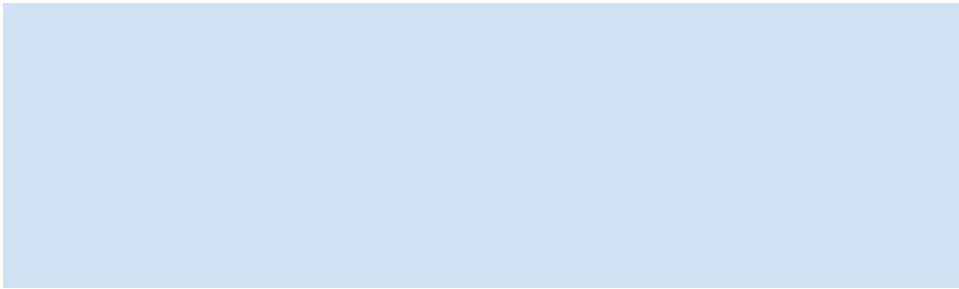
We get the CF:



Note the exponential decay factor.

# Finding CF with decoupling method

The equations of motion



$$q_1 = \frac{1}{\sqrt{2}}(x + y)$$

can be decoupled with the normal coordinates

$$q_2 = \frac{1}{\sqrt{2}}(x - y)$$

to yield the 2<sup>nd</sup> order homogeneous differential equations

$$\ddot{q}_1 + \frac{\gamma}{m}\dot{q}_1 + \omega_1^2 q_1 = 0 \quad \ddot{q}_2 + \frac{\gamma}{m}\dot{q}_2 + \omega_2^2 q_2 = 0 \quad \text{with}$$

$$\omega_1^2 = \frac{g}{l}$$

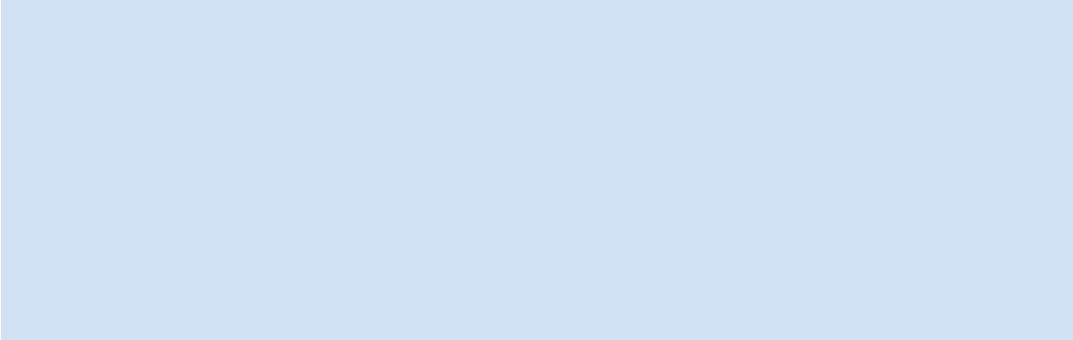
$$\omega_2^2 = \frac{g}{l} + 2\frac{k}{m}$$

that can be solved through trial solution  $q = \text{Re}(e^{i\omega t})$  to give same results

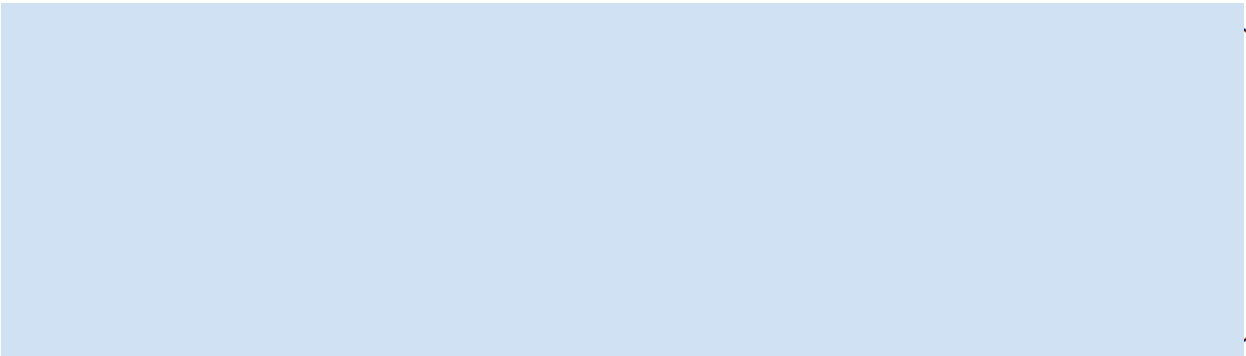


# Damped driven coupled pendula: finding PI

We have the CF. Now we need to find the PI, *i.e.* a solution to the full equation


$$\left. \right\} \begin{pmatrix} x \\ y \end{pmatrix} = \frac{F}{m} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \operatorname{Re}[\exp(i\alpha t)]$$

Try this ansatz  $\begin{pmatrix} x \\ y \end{pmatrix} = \operatorname{Re} \left[ \begin{pmatrix} P \\ Q \end{pmatrix} e^{i\alpha t} \right]$  which means solving the following


$$\left. \right\} \begin{pmatrix} P \\ Q \end{pmatrix} = \frac{F}{m} \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$


# Damped driven coupled pendula: finding PI

We have matrix equation of the sort:  $MU=V$

$$\overbrace{\left( \begin{array}{c} \dots \\ \dots \\ \dots \end{array} \right)}^M \begin{pmatrix} U \\ V \end{pmatrix} = \frac{F}{m} \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

and so  $U=M^{-1}V$ , i.e.  $\begin{pmatrix} P \\ Q \end{pmatrix} = M^{-1} \frac{F}{m} \begin{pmatrix} 1 \\ 0 \end{pmatrix}$

# Damped driven coupled pendula: finding PI

Need to find the inverse of  $M =$  

$$M^{-1} = \frac{1}{\det M} \text{adj}(M)$$

$$\begin{aligned} \det M &= \left[ -\alpha^2 + i\alpha \frac{\gamma}{m} + \left( \frac{g}{l} + \frac{k}{m} \right) \right]^2 - \left[ \frac{k}{m} \right]^2 \\ &= \left[ -\alpha^2 + i\alpha \frac{\gamma}{m} + \left( \frac{g}{l} + \frac{k}{m} \right) - \frac{k}{m} \right] \cdot \left[ -\alpha^2 + i\alpha \frac{\gamma}{m} + \left( \frac{g}{l} + \frac{k}{m} \right) + \frac{k}{m} \right] \\ &= \left[ -\alpha^2 + \frac{g}{l} + i\alpha \frac{\gamma}{m} \right] \cdot \left[ -\alpha^2 + \left( \frac{g}{l} + \frac{2k}{m} \right) + i\alpha \frac{\gamma}{m} \right] \\ &= B_1 e^{-i\theta_1} \cdot B_2 e^{-i\theta_2} \end{aligned}$$

$$\text{where } B_{1,2} = \left( (\omega_{1,2}^2 - \alpha^2)^2 + \left( \frac{\alpha\gamma}{m} \right)^2 \right)^{\frac{1}{2}} \text{ and } \tan \theta_{1,2} = \frac{-\alpha\gamma/m}{(\omega_{1,2}^2 - \alpha^2)}$$

# Damped driven coupled pendula: finding PI

Now  $\text{adj } M =$  

and we can write

$$\begin{aligned} -\alpha^2 + i\frac{\gamma}{m}\alpha + \frac{g}{l} + \frac{k}{m} &= \frac{1}{2}\left(-\alpha^2 + i\frac{\gamma}{m}\alpha + \frac{g}{l}\right) + \frac{1}{2}\left(-\alpha^2 + i\frac{\gamma}{m}\alpha + \frac{g}{l} + \frac{2k}{m}\right) \\ &= \frac{1}{2}\left[B_1 e^{-i\theta_1} + B_2 e^{-i\theta_2}\right] \\ \frac{k}{m} &= \frac{1}{2}\left(-\alpha^2 + i\frac{\gamma}{m}\alpha + \frac{g}{l} + \frac{2k}{m}\right) - \frac{1}{2}\left(-\alpha^2 + i\frac{\gamma}{m}\alpha + \frac{g}{l}\right) \\ &= \frac{1}{2}\left[B_2 e^{-i\theta_2} - B_1 e^{-i\theta_1}\right] \end{aligned}$$

and so

$$\text{adj } M = \frac{1}{2} \left( \img alt="A large light blue rectangular box representing a redacted matrix." data-bbox="355 800 795 935"/> \right)$$

# Damped driven coupled pendula: finding PI

$$M^{-1} = \frac{1}{\det M} \text{adj}(M) =$$

$$=$$

Recall 
$$\begin{pmatrix} x \\ y \end{pmatrix} = \frac{F}{m} \text{Re} \left[ M^{-1} \begin{pmatrix} 1 \\ 0 \end{pmatrix} e^{i\alpha t} \right]$$

→ 
$$\begin{pmatrix} x \\ y \end{pmatrix} = \dots$$
 we have it !

with 
$$B_{1,2} = \left( (\omega_{1,2}^2 - \alpha^2)^2 + \left( \frac{\alpha\gamma}{m} \right)^2 \right)^{\frac{1}{2}}$$
 and 
$$\tan \theta_{1,2} = \frac{-\alpha\gamma / m}{(\omega_{1,2}^2 - \alpha^2)}$$

# Finding PI with decoupling method

Inhomogeneous equations written in terms of normal coordinates

$$\ddot{q}_1 + \frac{\gamma}{m} \dot{q}_1 + \omega_1^2 q_1 = \frac{F}{m} \cos \alpha t \quad (1)$$

$$\ddot{q}_2 + \frac{\gamma}{m} \dot{q}_2 + \omega_2^2 q_2 = \frac{F}{m} \cos \alpha t \quad (2)$$

$$q_1 = \frac{1}{\sqrt{2}}(x + y) \quad q_2 = \frac{1}{\sqrt{2}}(x - y)$$

Trial ansatz for (1)  $q_1 = \text{Re}[A_1 \exp(i\alpha t)]$

$$\Rightarrow \left( -\alpha^2 + i\alpha \frac{\gamma}{m} + \frac{g}{l} \right) A_1 = \frac{F}{m} \quad \text{with}$$

$$A_1 = \frac{(F/m) \exp(i\theta_1)}{\left( (\omega_1^2 - \alpha^2)^2 + (\alpha\gamma/m)^2 \right)^{1/2}}$$

$$\tan \theta_1 = \frac{-(\alpha\gamma/m)}{(\omega_1^2 - \alpha^2)}$$

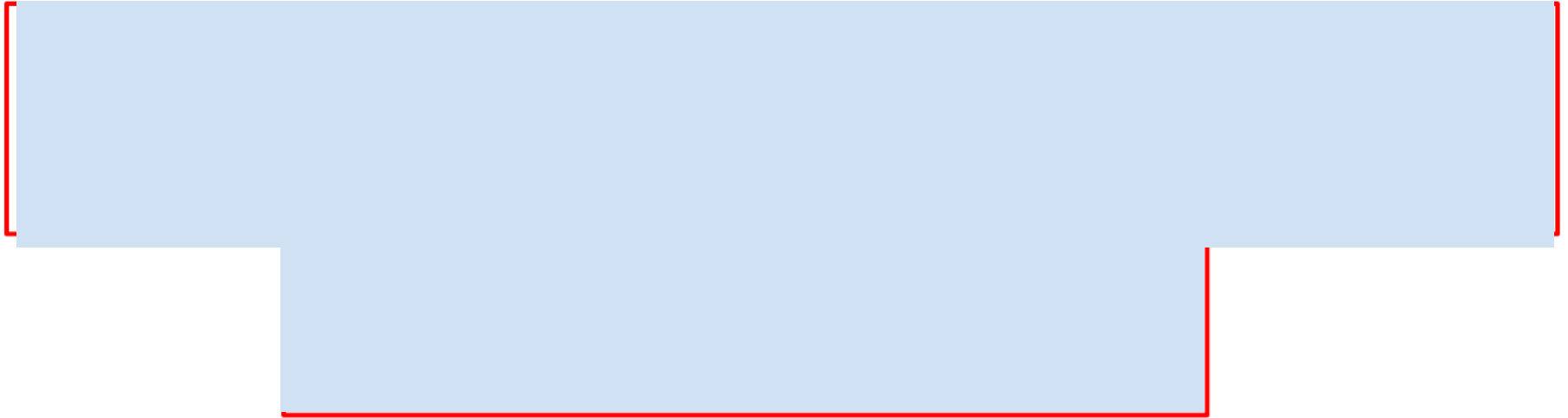
Hence



Same procedure for (2) gives entirely analogous expression for  $q_2$ .  
From these same expressions are obtained for  $x$  and  $y$  as before.

# Damped driven coupled pendula: full solution

Solution = CF + PI

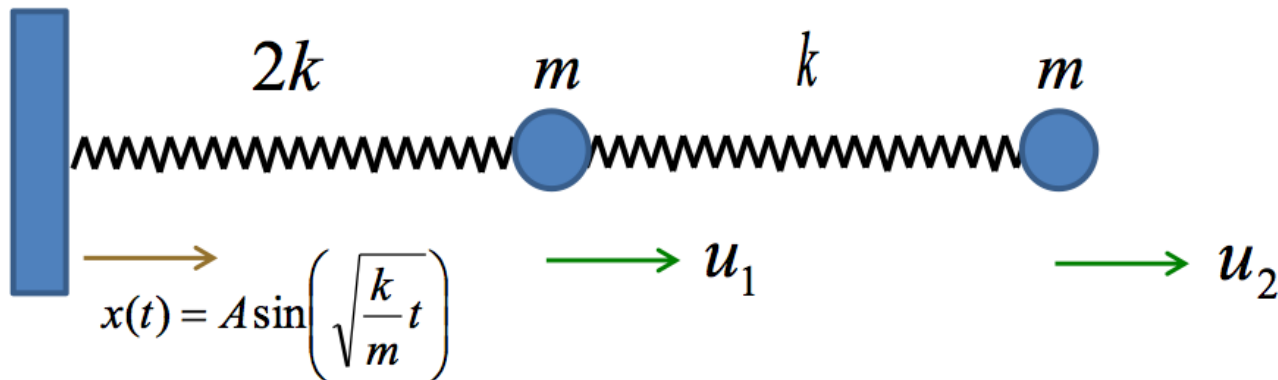


$$\text{with } B_{1,2} = \left( (\omega_{1,2}^2 - \alpha^2)^2 + \left( \frac{\alpha\gamma}{m} \right)^2 \right)^{\frac{1}{2}}, \quad \tan \theta_{1,2} = \frac{-\alpha\gamma/m}{(\omega_{1,2}^2 - \alpha^2)}, \quad \omega_1^2 = \frac{g}{l} \quad \& \quad \omega_2^2 = \frac{g}{l} + 2\frac{k}{m}$$

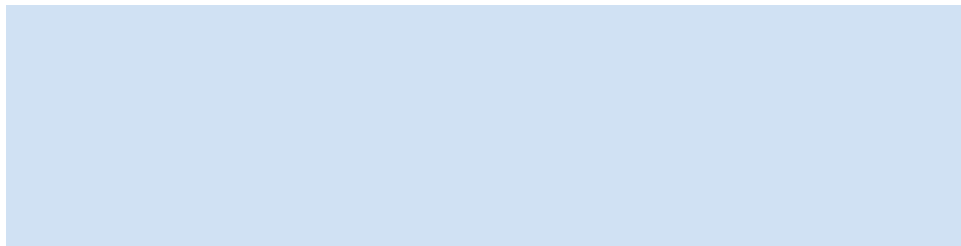
The CF part is the 'transient solution' determined by the initial conditions;  
the PI part is the 'steady state solution' determined by the driving force.

# Horizontal spring-mass system with driving term

Consider two masses moving without friction, with two springs of spring constants  $2k$  and  $k$  respectively, connected to wall which is driven by an external force to have time-dependent displacement  $x(t) = A \sin\left(\sqrt{\frac{k}{m}}t\right)$



Equations of motion:





# Horizontal spring-mass system with driving term – find the CF

Write down the homogeneous case  
and find CF using matrix method

Try  $\begin{pmatrix} u_1 \\ u_2 \end{pmatrix} = \text{Re} \left( \begin{pmatrix} U_1 \\ U_2 \end{pmatrix} e^{i\omega t} \right)$  gives  $\begin{pmatrix} \phantom{U_1} \\ \phantom{U_2} \end{pmatrix} = 0$

Requiring determinant = 0 yields

Substitute back in to eigenvector equation  $\begin{pmatrix} U_2 \\ U_1 \end{pmatrix}_1 = 1 - \sqrt{2}$  and  $\begin{pmatrix} U_2 \\ U_1 \end{pmatrix}_2 = 1 + \sqrt{2}$

# Horizontal spring-mass system with driving term – find the PI

We have the CF. Now we need to find the PI, *i.e.* a solution to the full equation

$$\left( \begin{array}{c} \text{[Redacted]} \\ \text{[Redacted]} \end{array} \right) = \frac{Ak}{m} \begin{pmatrix} 2 \\ 0 \end{pmatrix} \text{Re}[\exp(i\sqrt{\frac{k}{m}}t)]$$

Try ansatz  which means solving

$$\begin{pmatrix} \frac{2k}{m} & -\frac{k}{m} \\ -\frac{k}{m} & 0 \end{pmatrix} \begin{pmatrix} P \\ Q \end{pmatrix} = \frac{kA}{m} \begin{pmatrix} 2 \\ 0 \end{pmatrix}$$

Inverse of LHS 2x2 matrix is

$$\frac{1}{-(k/m)^2} \begin{pmatrix} 0 & \frac{k}{m} \\ \frac{k}{m} & \frac{2k}{m} \end{pmatrix} = \frac{m}{k} \begin{pmatrix} 0 & -1 \\ -1 & -2 \end{pmatrix}$$

and so

# Horizontal spring-mass system with driving term – full solution

Solution = CF + PI

