



# Deterministic Entanglement of Trapped-Ion Spin-Qubits

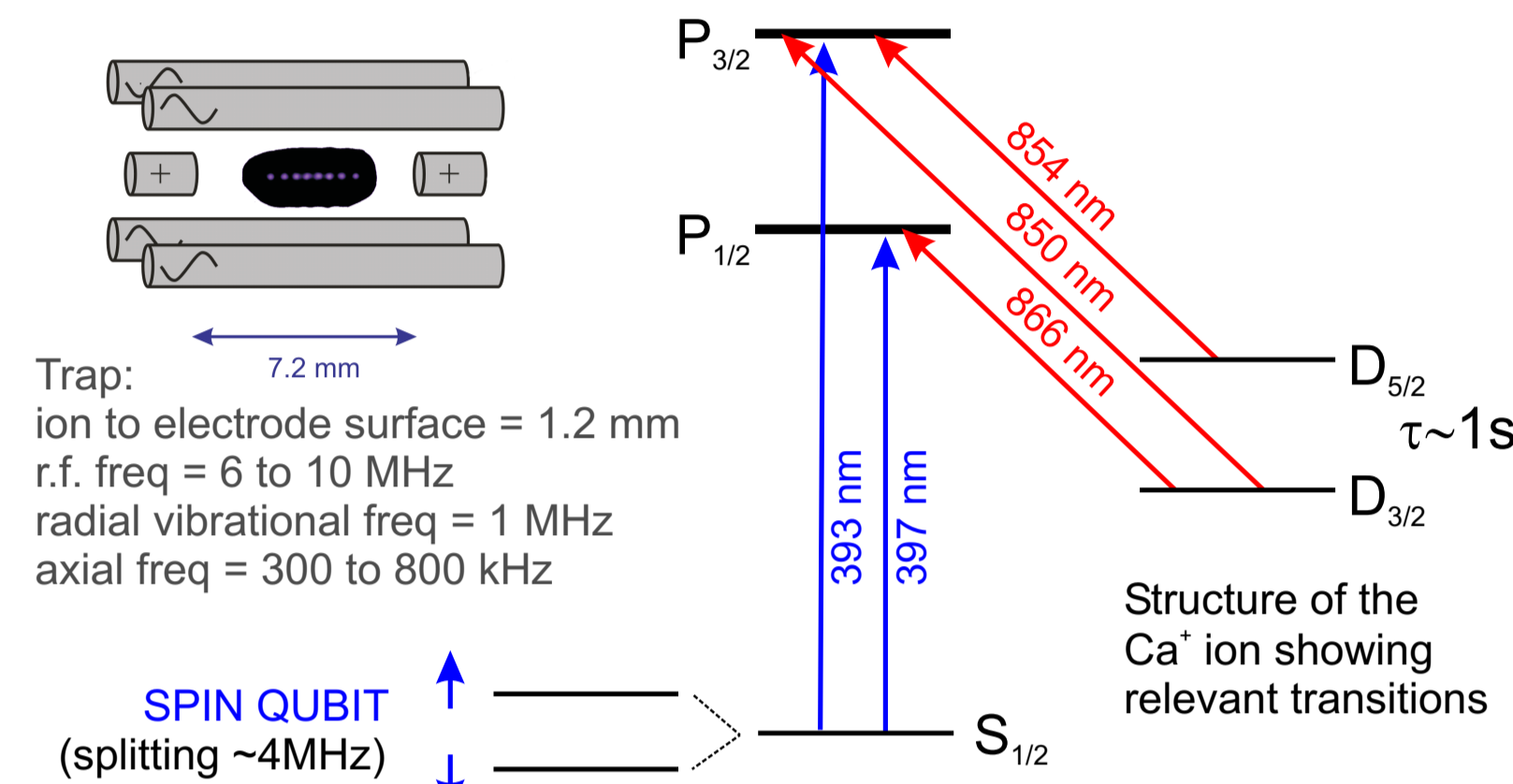
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## Main points

We present experiments and theory in quantum information processing using trapped ions.  
This poster concentrates on entanglement and gates: see accompanying poster for cooling, coherence.



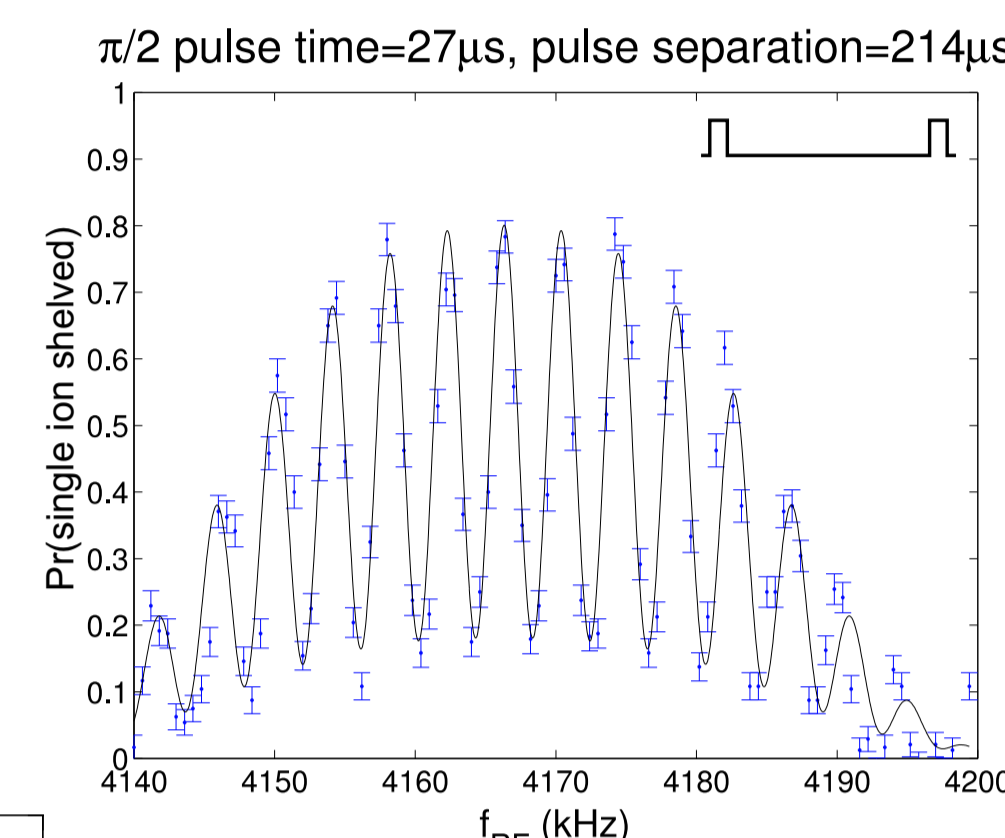
- Summary of Results**
- 10 two-ion (2 qubit) Rabi flops with high visibility
  - Deterministic entanglement of 2 ions (calcium 40 spin qubits) at **82(2)% fidelity**
  - Schrodinger cat with 1 ion and motion:
    - $\alpha$  up to 3.5(3) ( $\langle n \rangle = 12$ )
    - well outside Lamb-Dicke regime:  $\eta^2 2n = 1.6$
    - $\alpha = 1$  preserved for 422  $\mu$ s with 80(20)% fidelity
    - also  $\alpha = -2, 0, +2$  with 2 ions
  - robust convenient tomography method
  - (th.) factorization of general phase gates (ask for details)
  - (th.) composite pulses for fast gate ( $t=1/\text{trap freq}$ ) insensitive to optical phase

## Single-qubit gates, 1-2 ions

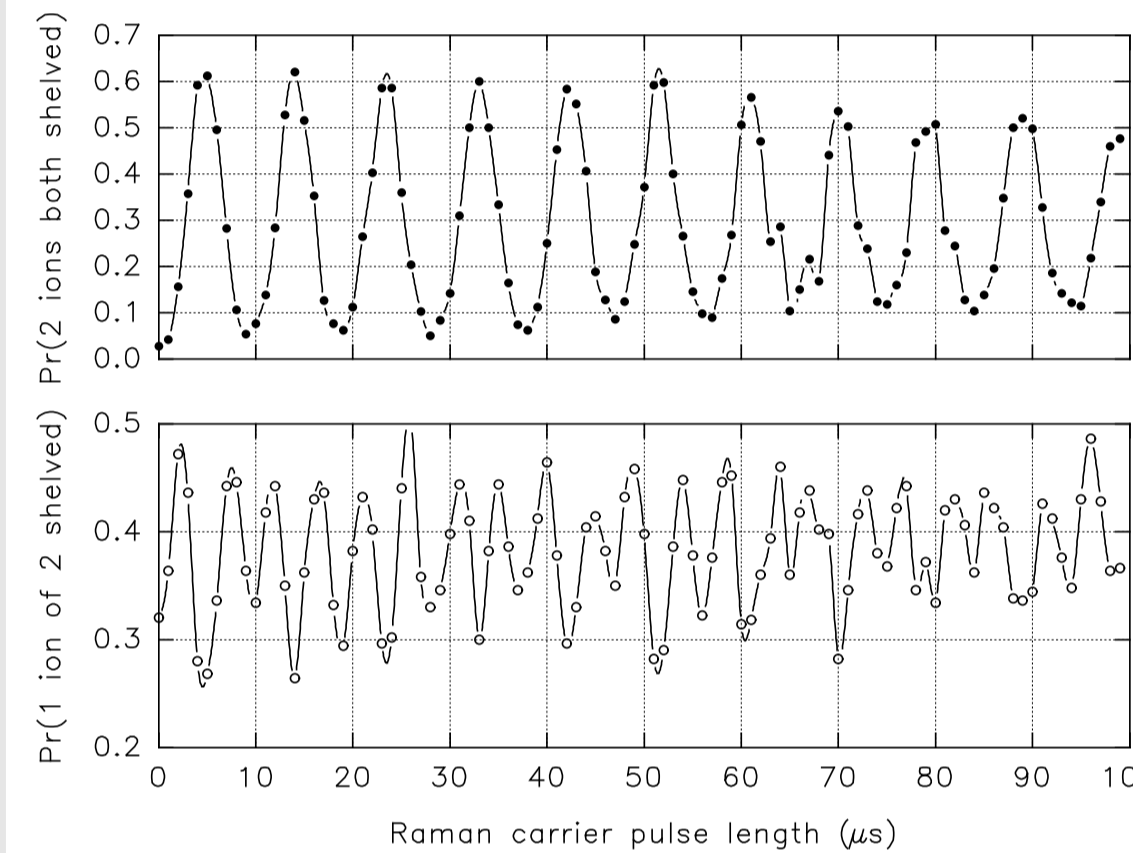
Spin qubit state coherently manipulated either by magnetic resonance or by stimulated Raman transition.

### Single-ion Ramsey fringes

This data is for a two-pulse Ramsey sequence using magnetic resonance with a single trapped ion. Interference fringes are seen as the RF frequency is scanned.



### Two-ion Rabi oscillations



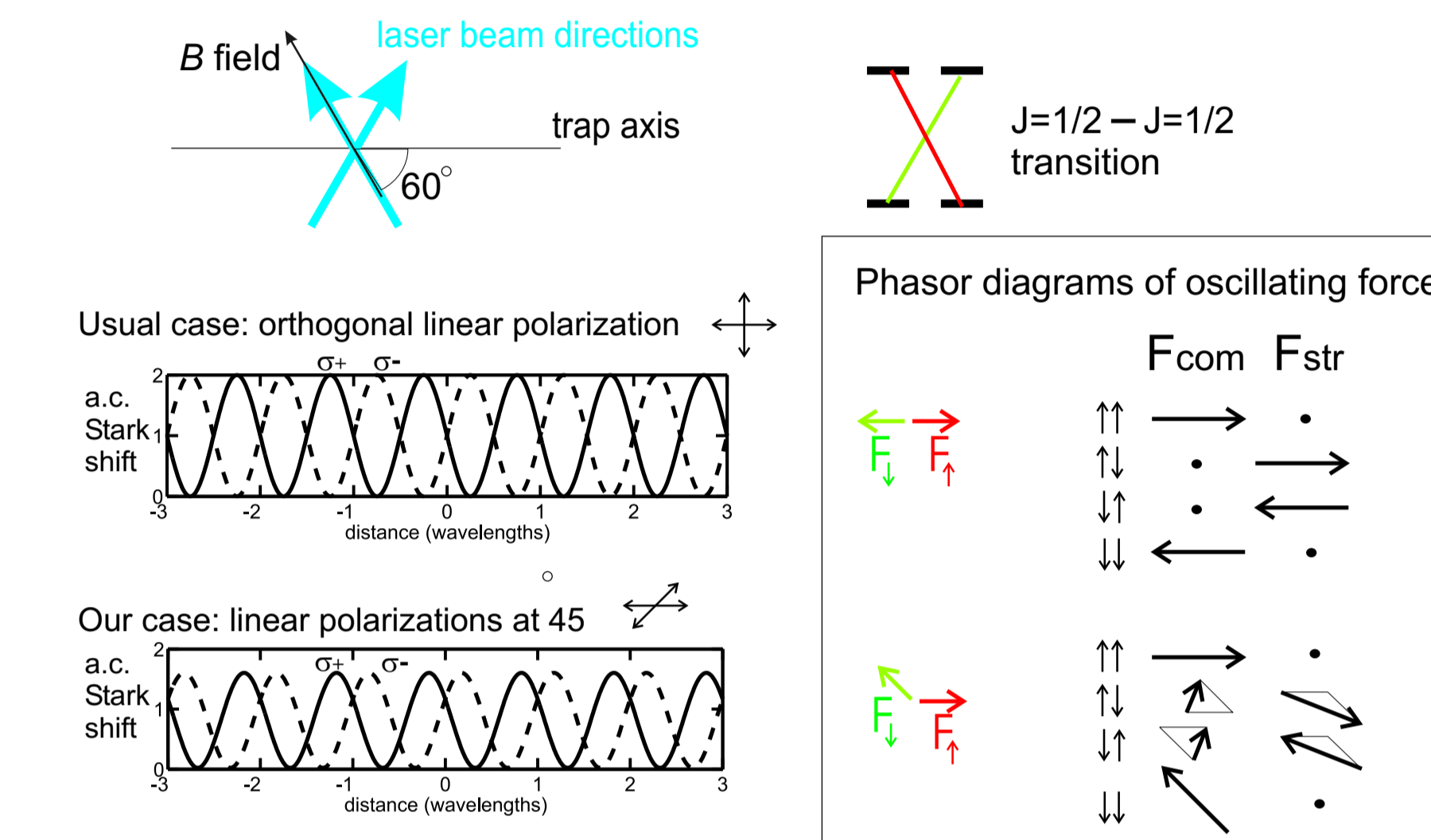
2-ion spin-state = 2 qubits. Rabi flopping, here driven by the Raman transition (4.5  $\mu$ s  $\pi$ -time), gives a single-qubit rotation applied to both qubits simultaneously.

EIT readout  $\rightarrow 0, 1$  or 2 ions shelved  
 $\rightarrow$  infer  $P(\uparrow\uparrow)$ ,  $P(\uparrow\downarrow) + P(\downarrow\uparrow)$ ,  $P(\downarrow\downarrow)$

Coherence time (measured separately) of order 1 ms: the falling visibility here is a beating effect owing to unequal illumination of the ions.

## Spin-dependent force

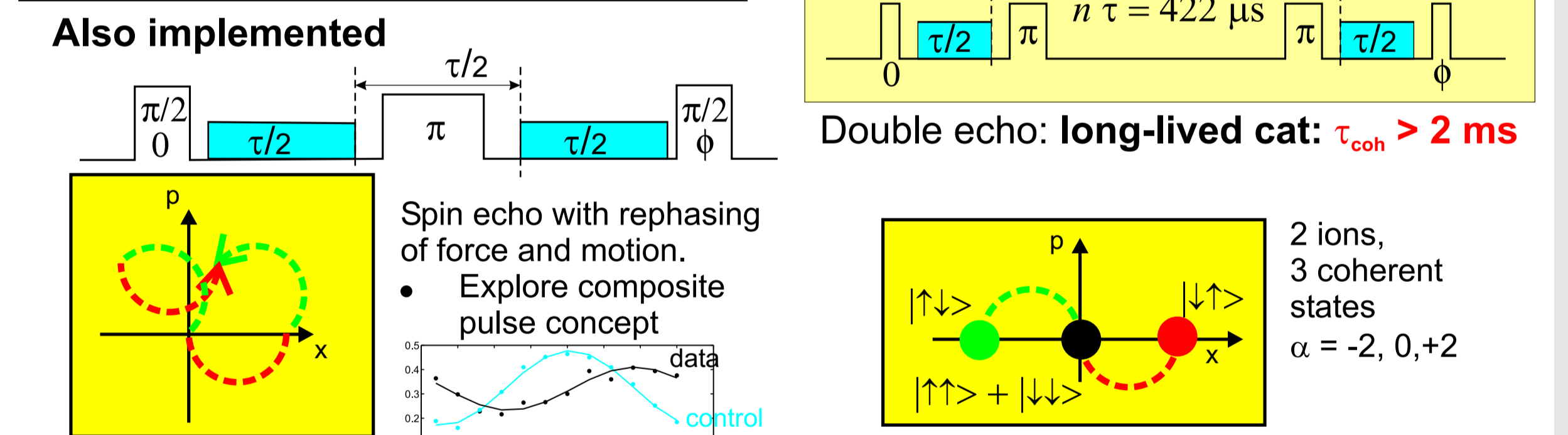
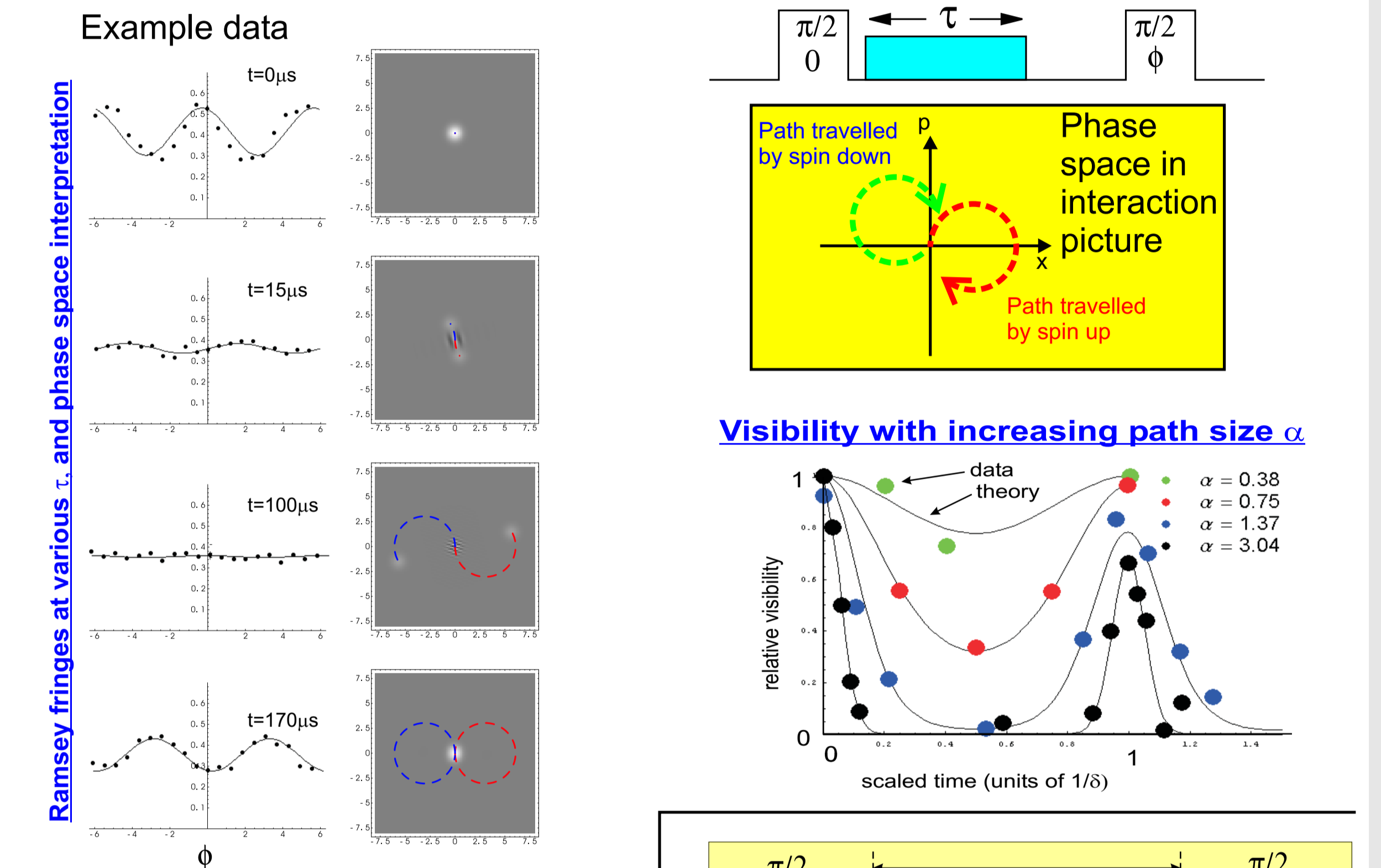
For two-qubit gates we use spin-dependent forces: push ions depending on spin state  
 $\rightarrow$  Coulomb interaction gives a two-qubit phase.  
The force is an optical dipole force in a standing wave with polarization gradient.



- A classical force displaces the motional state in phase space. e.g. oscillating force drives the state around a loop.
- Initial  $|n=0\rangle$  displaced  $\rightarrow$  Glauber coherent state
- When loop closes one has  $U = \text{diag}(\exp(i\phi))$
- We have found a general mathematical method to decompose such U for multiple ions as a unique product of 1, 2, ... n-qubit phase gates.**

## Schrödinger Cat experiments

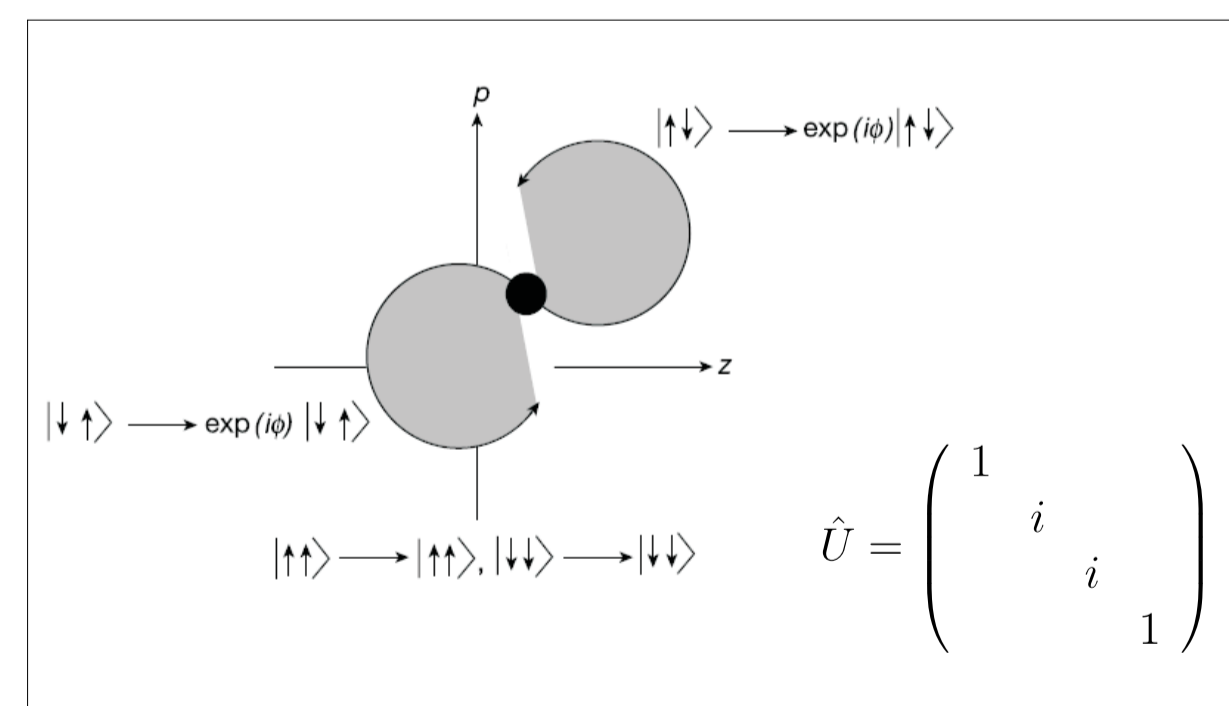
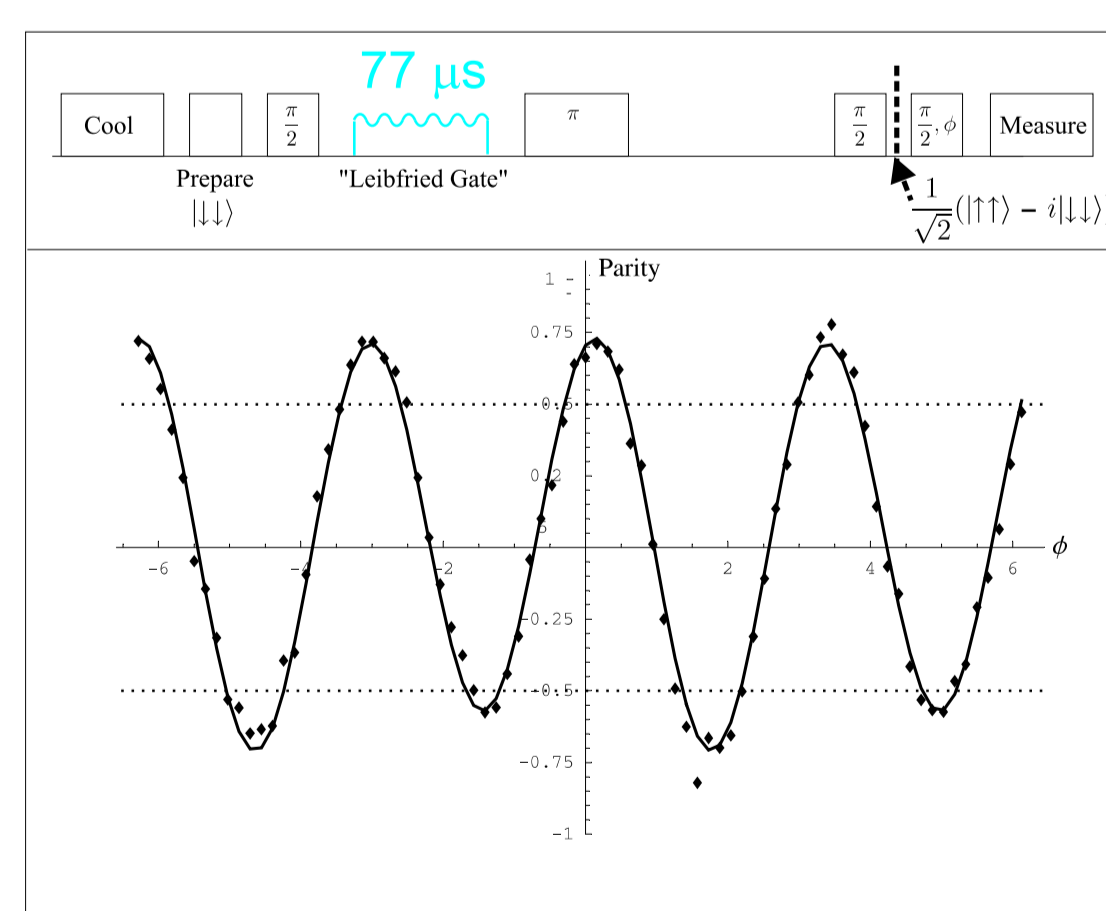
Coherent states of a harmonic oscillator approximate to classical motion, and a superposition of such states at mesoscopic excitation  $\langle n \rangle$  is a type of Schrödinger cat. Oscillating spin-dependent force  $\rightarrow$  create such mesoscopic superpositions with single or pairs of ions. Spin state = measuring device entangled with the motion. We prove the 'cat' maintains its coherence by bringing the two parts back together and observing an interference. [As first demonstrated by Monroe *et al.* Science 272 1131 (1996).]



## Deterministic entanglement

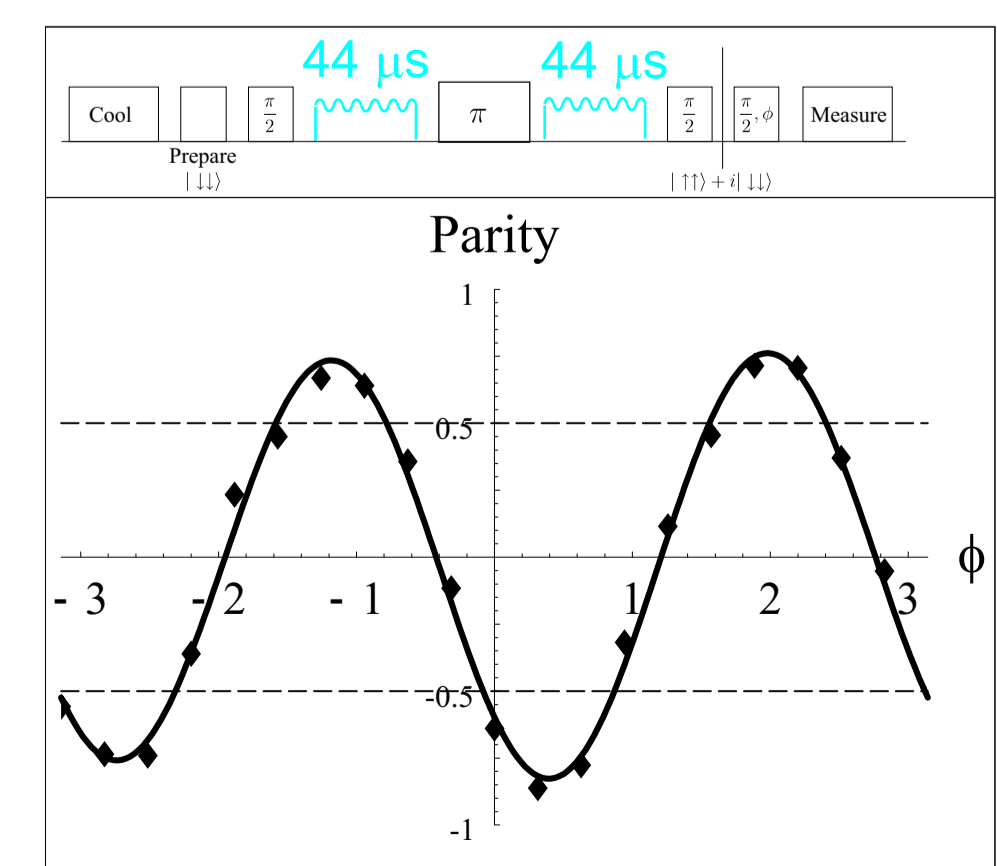
- Deterministic** (i.e. single-shot, no post-selection) entanglement of 2 spin-qubits
- gate uses same oscillating spin-dependent driving force as to create Schrödinger cats, with force frequency close to  $\omega_{str}$  & ion separation = integral number of standing wave periods
- $\Rightarrow$  only stretch mode excited
- $\Rightarrow$  states  $\uparrow\downarrow, \downarrow\uparrow$  acquire a phase;  $\uparrow\uparrow, \downarrow\downarrow$  do not.

(Leibfried *et al.* [Nature 422 412 (2003)])



### Results

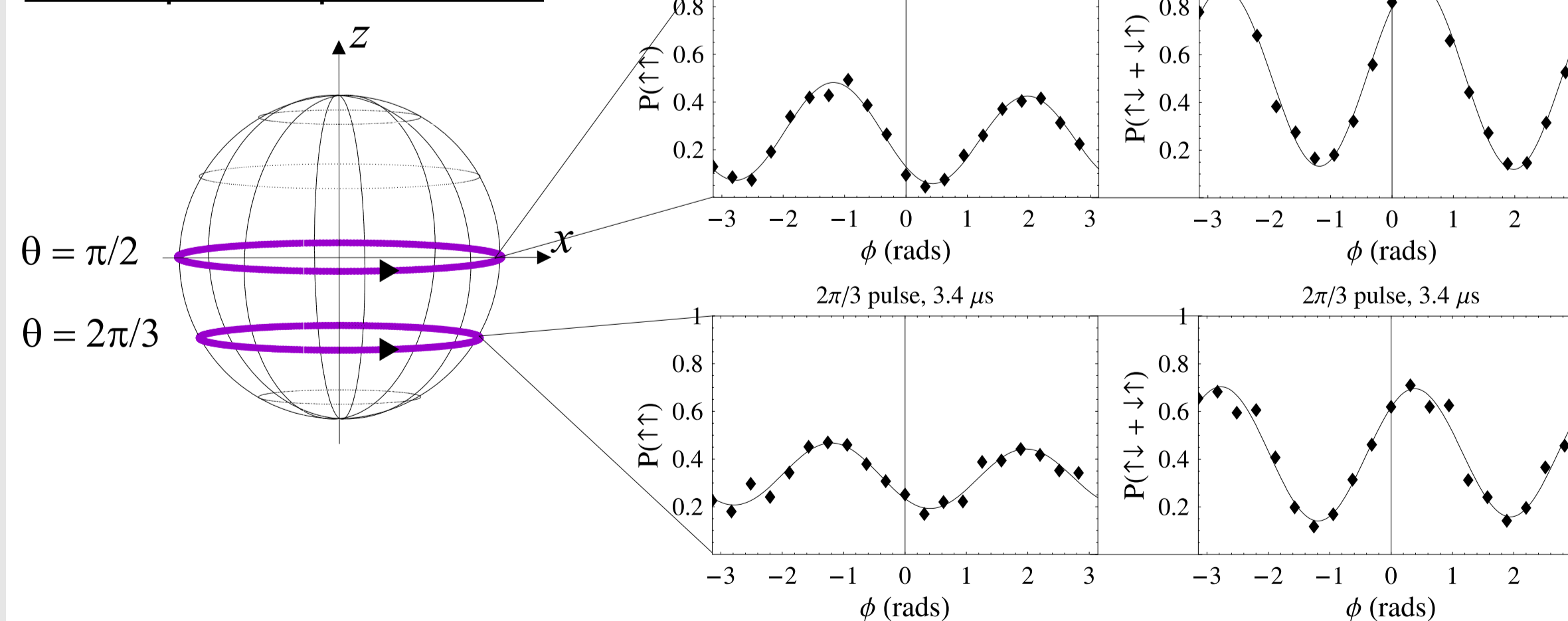
- Spin echo sequence to suppress slow drift effects
- $V_{com} = 500$  kHz,
- ion sep =  $9 \mu$ m =  $22 \lambda$
- the entangled state ( $\uparrow\uparrow - i\downarrow\downarrow$ ) is produced
- a further  $\pi/2$  analysis pulse with variable phase  $\phi$  demonstrates  $\cos(2\phi)$  oscillations in the parity signal with amplitude  $> 0.5$
- 1st exp: parity amplitude  $\Rightarrow$  entangled state fidelity  $> 75(5)\%$
- 2nd exp: two loops, one in each half of spin-echo: fidelity 82(2)%



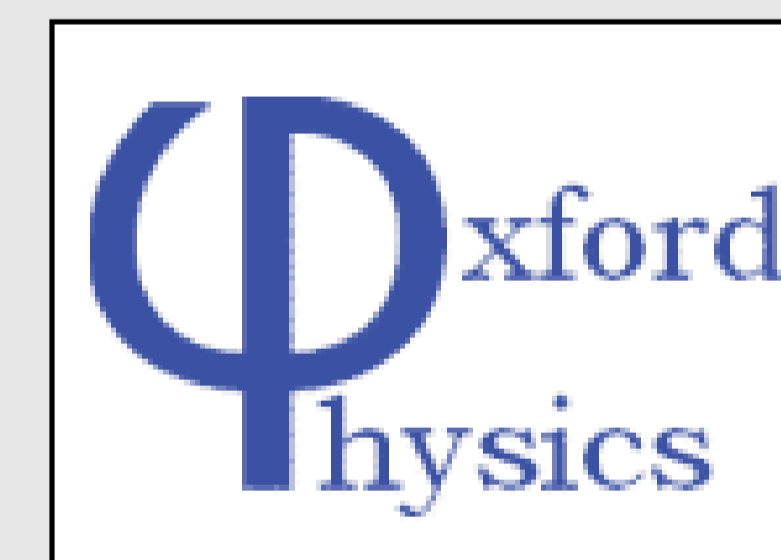
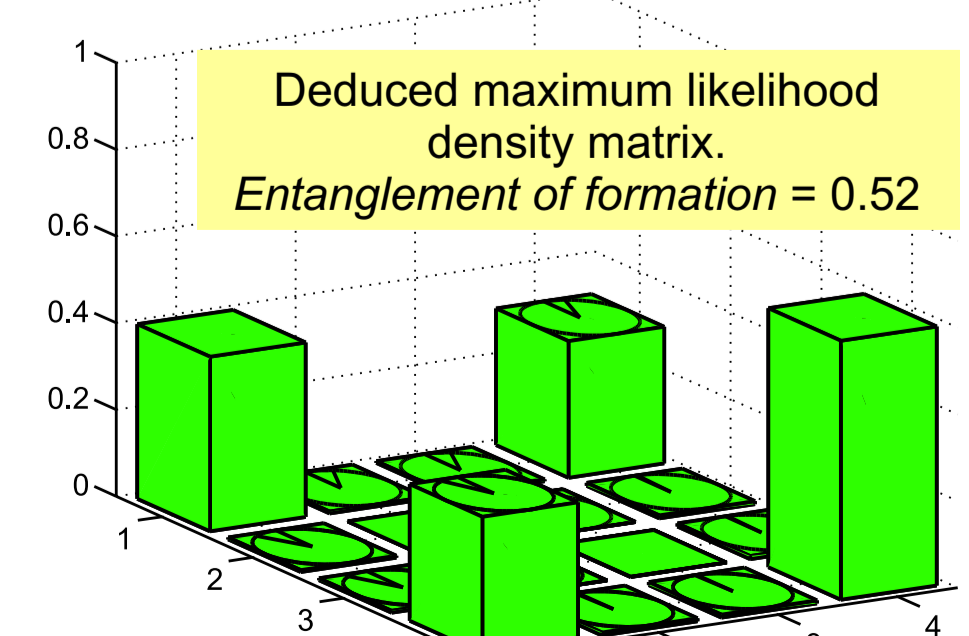
## Tomography

In general, tomography involves accumulating information by applying well-chosen rotations to the qubits and measuring them in a fixed basis. We developed a convenient scheme which is robust against typical experimental issues. The rotation (same for all qubits) is through  $\theta$  about an axis  $\phi$  in the x-y plane.  $\phi$  is scanned from  $-\pi$  to  $\pi$ :  $P(\text{spin state}) = \sum(\text{sinusoidal functions of } \phi)$ : this allows robust curve-fitting of sin functions with period  $\pi$  and  $2\pi$  and an offset. Each contribution to the fit yields 1 or 2 real numbers; two values of  $\theta$  are needed for complete information. A maximum likelihood estimation method is then used to obtain the physical density matrix closest to that obtained from the data.

### Bloch Sphere representation



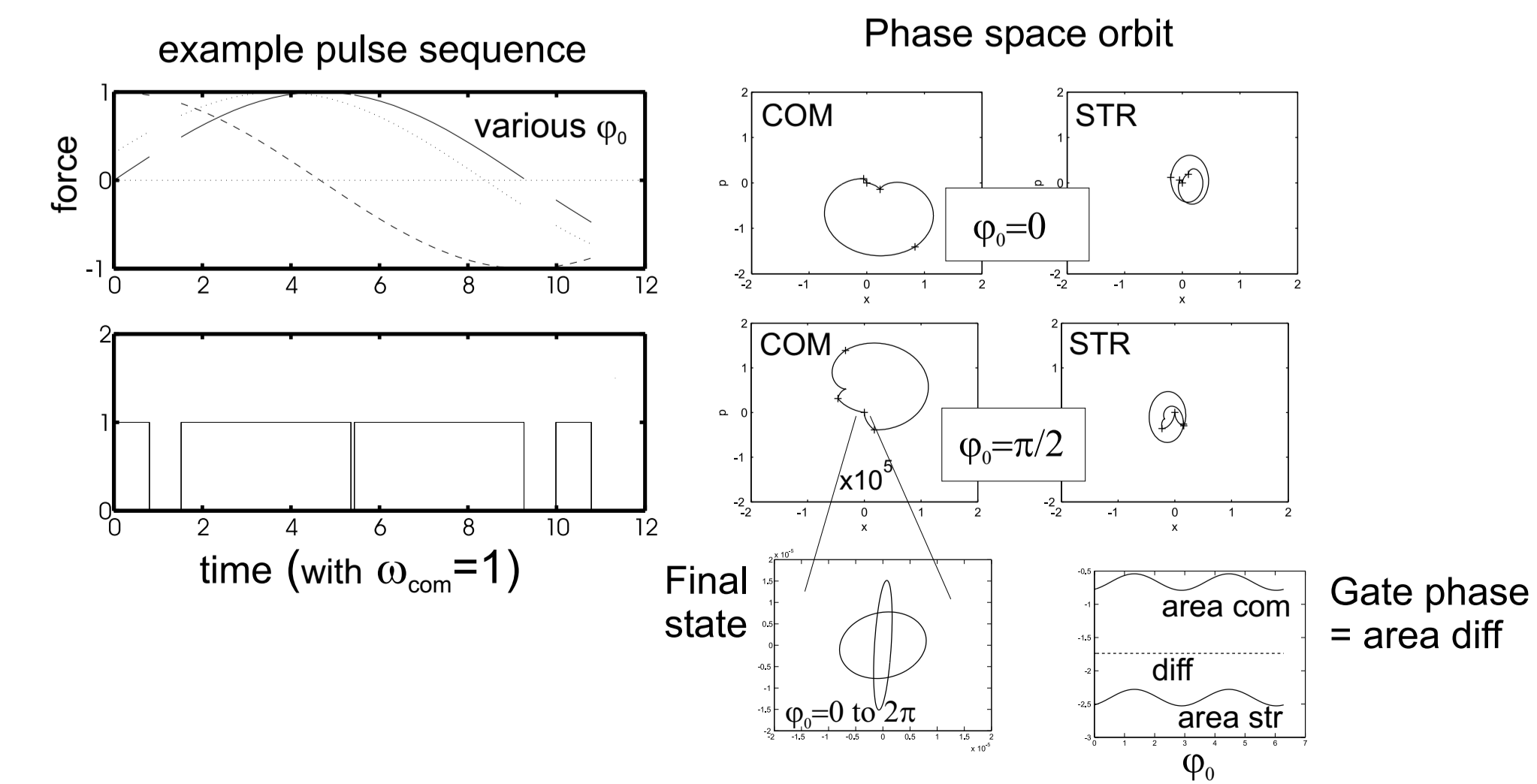
We show the results of tomography experiments on an entangled state  $|\uparrow\uparrow\rangle + e^{i\phi}|\downarrow\downarrow\rangle$ . The large amplitude of the oscillations with period  $\pi$  compared to those with period  $2\pi$  indicate that the coherence between  $\uparrow\uparrow$  and  $\downarrow\downarrow$  is large. (Our measurement method does not distinguish  $\uparrow\downarrow$  from  $\downarrow\uparrow$ , but this has little influence for this example).



## Composite pulses for fast robust gates

Wobble gate works well at low  $\delta = \omega - \omega_{str}$  but is slow,  $\tau = 2\pi/\delta$ . At high  $\delta$  both COM and STR modes excited, can't close both loops in a single pulse (incommensurate freq). Tailor  $f(t)$  in order to go faster? : lose insensitivity to optical phase  $\phi_0$ .

- We find fast composite pulse sequences maintaining insensitivity to  $\phi_0$
- Issues: loop closure, constant area, lightshift phase



top hat

$$f(t) = \sum_{n=1}^N T((t-t_n)/\tau_n) f_n \sin(\omega_n t + \phi_n)$$

random fixed

$$\phi_n = \phi_0 + \Delta\phi_n$$

$$\alpha(t) - \alpha(0) = \frac{i}{m\omega_0 a} \int_0^t e^{i\omega_0 t} f(t) dt$$

