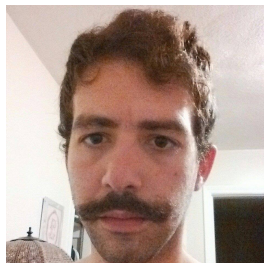


Black Hole Collapse in Large-C CFT

Julian Sonner



Oxford
13 July 2016



Tarek Anous (MIT)



Thomas Hartman (Cornell)



Antonin Rovai (Geneva)



SwissMAP

The Mathematics of Physics
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Introduction

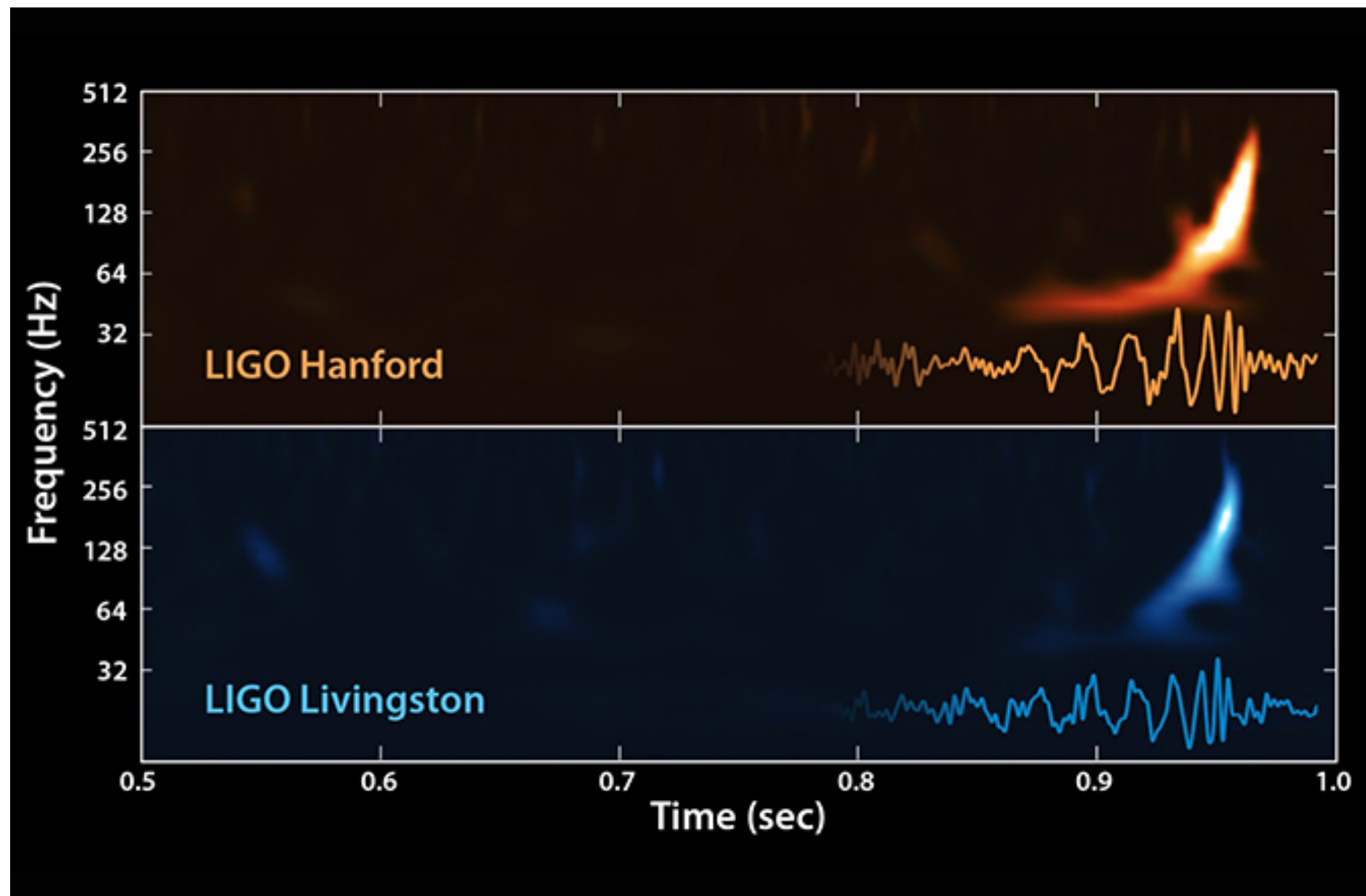


The black hole as movie star: Gargantua



‘Gargantua’, C. Nolan & K. Thorne

The black hole as news sensation: GW150914



‘Ligo Black Hole Binary’ R. Drever, R. Weiss, K. Thorne et mult. al.

$$S = \frac{k_B c^3 A}{4G_N \hbar}$$

black holes are **thermodynamic** systems

their entropy is proportional to the **area** of the event horizon



information loss paradox:

a BH formed from a pure state will evolve into a **mixed** state (of Hawking radiation)

holography:

a theory of quantum gravity should have **information** ~ **area**



General Plan

AdS/CFT relates gravity (often in AdS) to **unitary** field theory (often CFT)

Lots of progress **gravity** → **CFT** (my favorite: AdS/CMT)

Less known about **CFT** → **(quantum) gravity**

→ despite developments in CFT, CMT:

- time evolution and spread of entanglement
- **thermalization** of closed quantum systems (e.g. via eigenstates)
- non-perturbative methods (e.g. bootstrap)

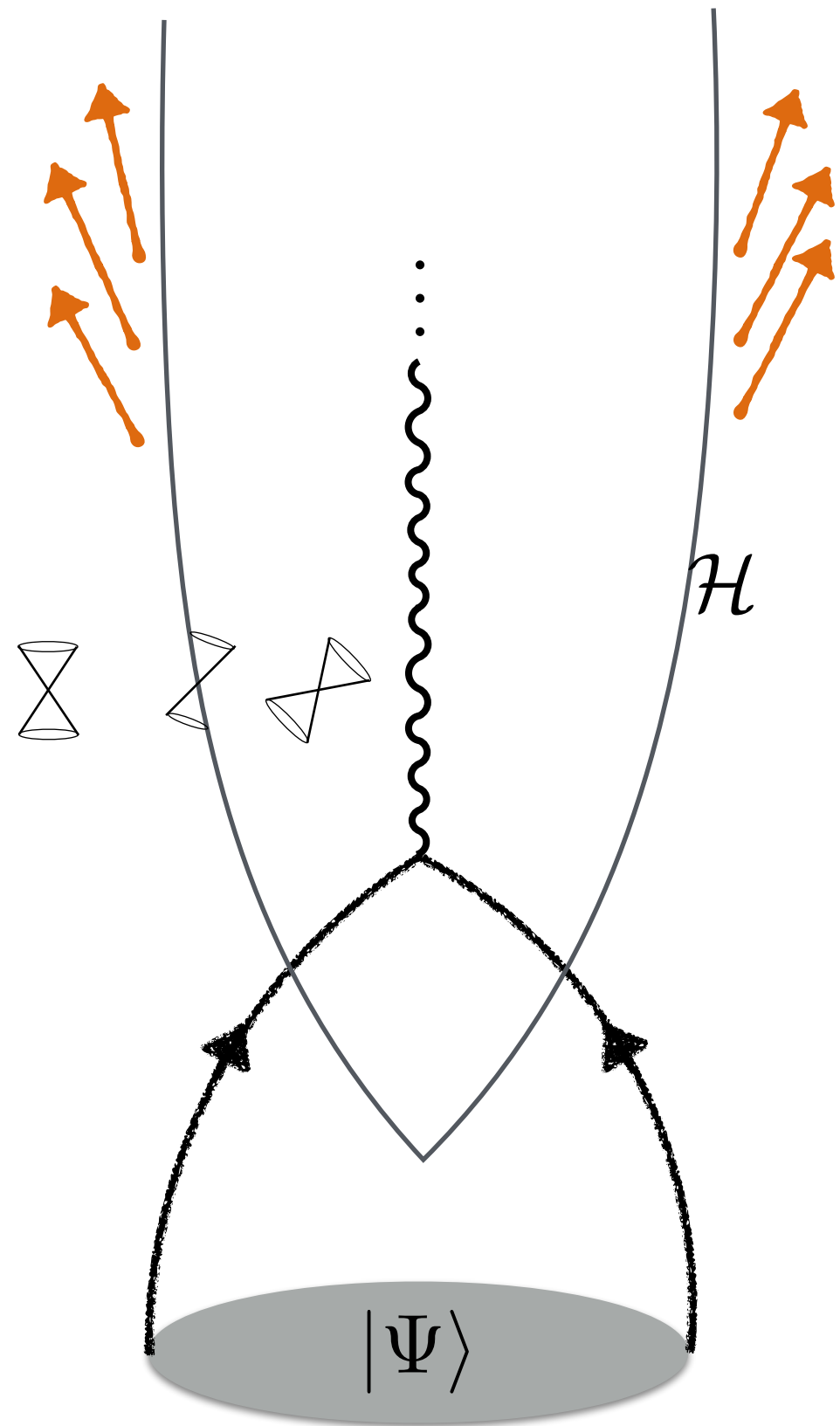
Thermalization → BH formation (& evaporation)

Outline

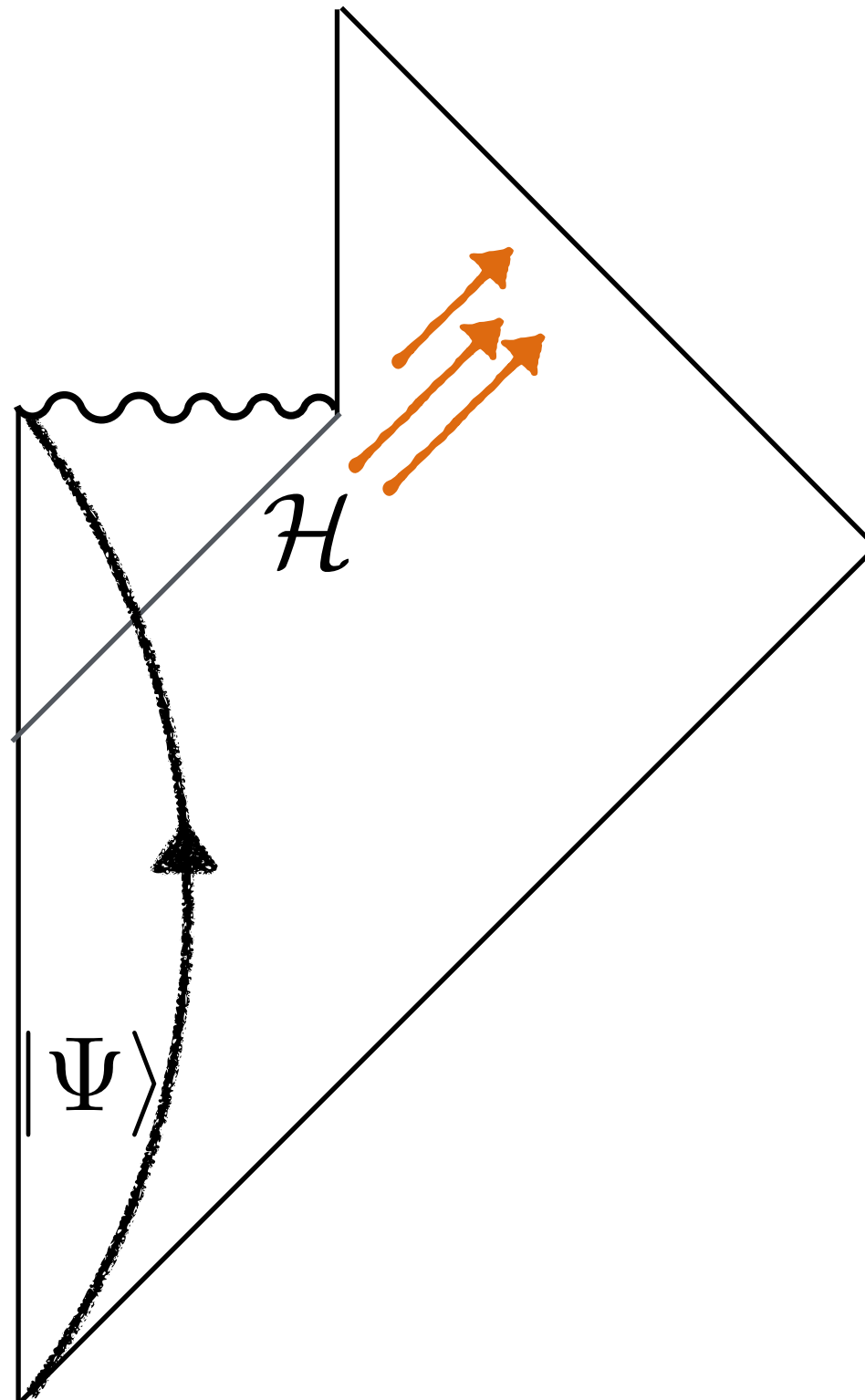
1. Reminder about black holes
2. CFT implosions
3. Thermalization @ large c
4. Conclusions

“the trouble with black holes”

- outgoing **Hawking radiation** is thermal $\rho_{\text{Gibbs}}(T_H)$
- a **horizon** \mathcal{H} cloaks the singularity
- initial pure state $|\Psi\rangle$ of matter collapses inwards

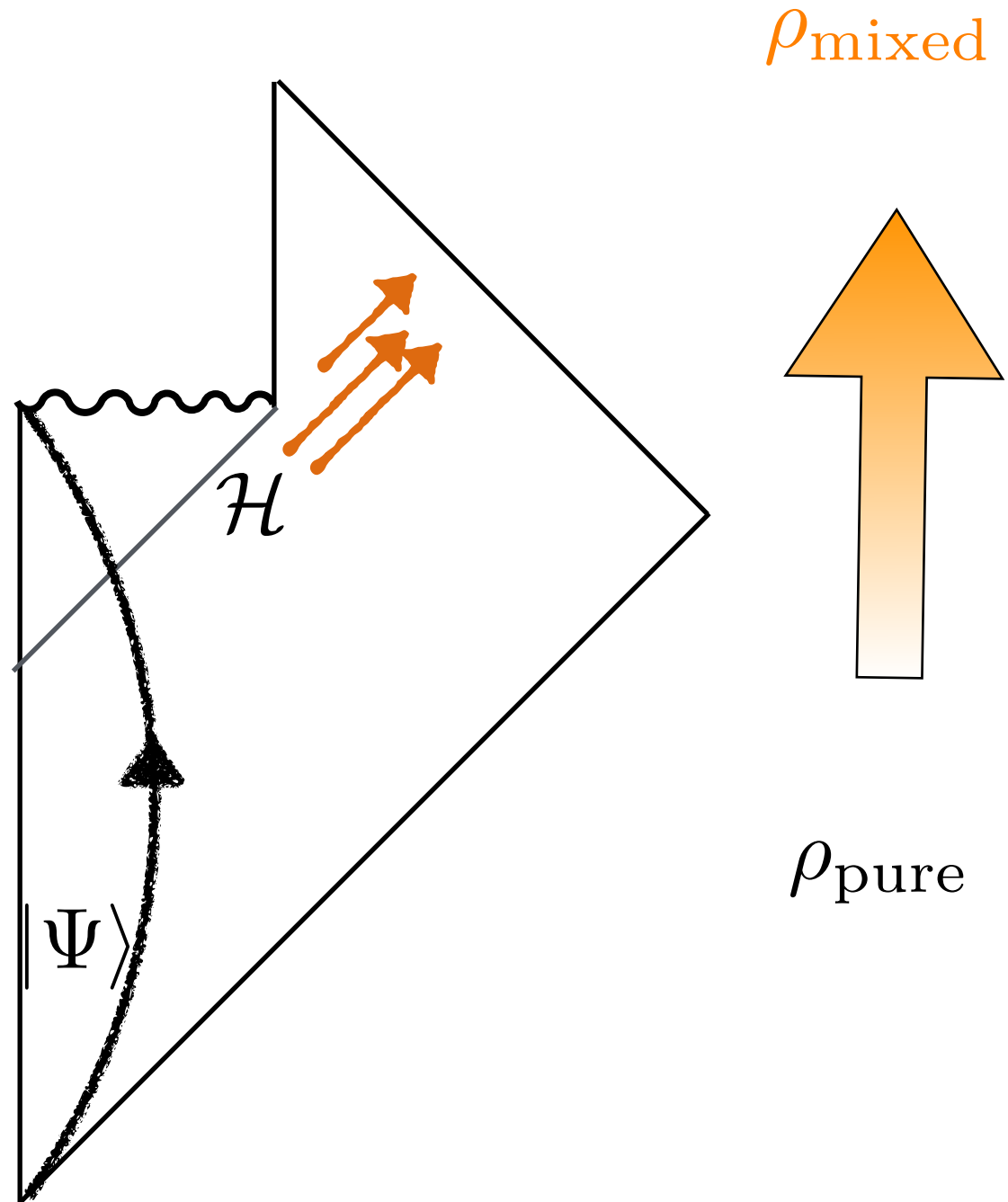


black holes evaporate



The Paradox

- gravity as an EFT implies pure to **mixed** evolution
 - fundamentally incompatible with a unitary S-matrix
1. quantum gravity is non-unitary
 2. gravity EFT makes no sense
 3. (subtle) corrections to Hawking result



Holography:

quantum gravity $=$ quantum field theory

hence AdS/CFT only allows for options 2 & 3.

the anti-information loss paradox:

how does an obviously unitary theory lose information?

The Plan

(of a first-principles calculation in holographic CFT)

1. define an initial state in CFT which forms a black hole
→ “quantum quench”
2. understand time evolution in strong-coupling regime
→ non-equilibrium CFT (monodromy method / CFT_2)
3. analyze suitable observables (e.g. correlations & EE)
→ strongly-coupled CFT thermalization
4. diagnose signs of information loss & recovery
→ unitarity constraints on correlations

Unitarity vs Thermalization

(constraints on long-time correlations from unitarity)

Correlations in a closed quantum system, e.g.

$$G(t) = \text{tr} \rho \mathcal{O}(t) \mathcal{O}(0)$$

Time average over a large time T **cannot vanish** by unitarity

$$\lim_{T \rightarrow \infty} \overline{|G(t)|^2} \neq 0$$

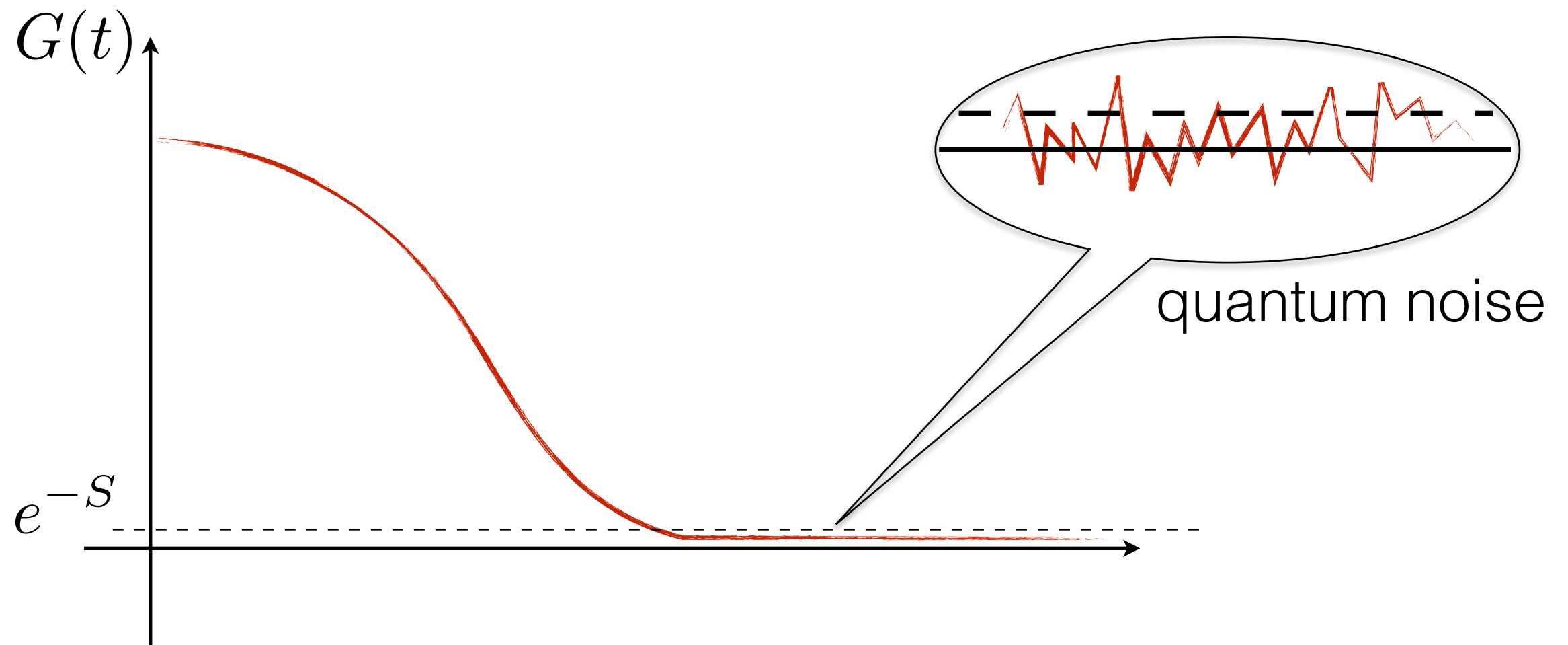
Need to assume spectrum is generic (no specific ordering principle)

→ can appeal e.g. to ETH so estimate $G(t) \sim e^{-S}$ at late times

Unitarity vs Thermalization

(constraints on long-time correlations from unitarity)

$$\rho = e^{-\beta H}$$



CFT implosions



Looking for the Right Place

0D matrix models(IOP,...): connection to geometry?

2D black hole (CGHS): solvable but very different

3D story shares salient features of 4D (and higher)
in fact central to micro-state counting success (D1-D5)

the trouble: no local degrees of freedom (Achucarro & Townsend):

$$S_{3D} = S_{CS}[A] - S_{CS}[\bar{A}]$$

other side of the coin: CFT₂ puts powerful tools at our disposal

3D Gravity + Matter

→ add matter: get local dof. BUT need new tools

focus on a universal sector, by defining a $1/c$ expansion:

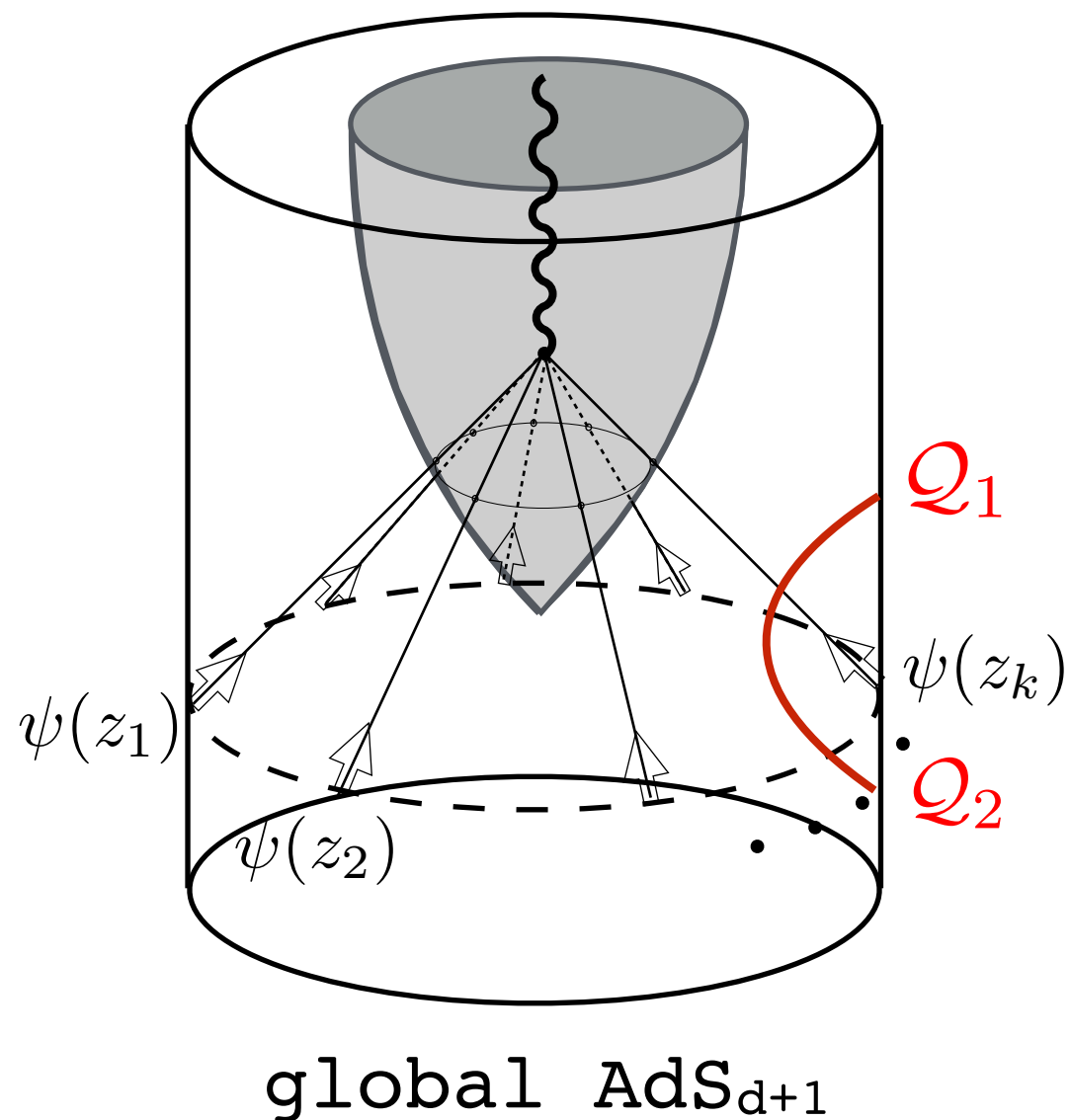
→ any microscopic theory in this class defines some 3D quantum gravity theory (sparse spectrum)

*3D gravity + matter non-trivial, but solvable
→ ideal place to study BH puzzles!*

From bulk point of view this is G_N expansion

CFT_2 gives a non-perturbative definition of quantum gravity.

The Black Hole in the Tin Can



Throw in a shell of n dust particles

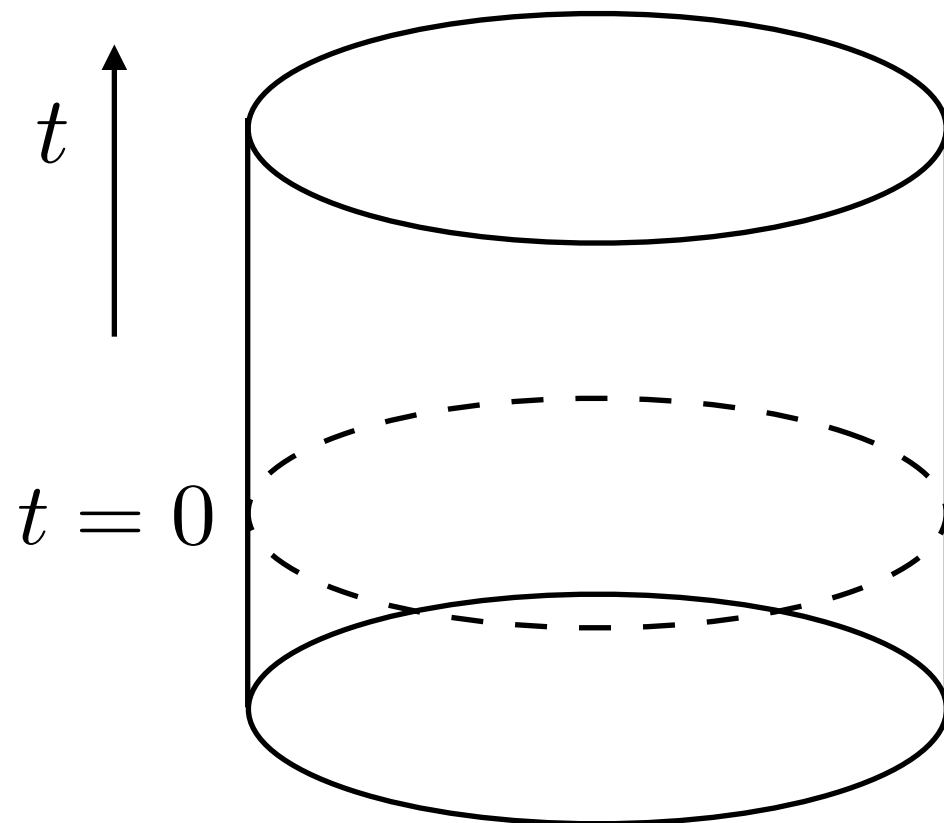
smooth limit: $n \rightarrow \infty$

BH collapse: Vaidya metric

Use light operators \mathcal{Q} to probe geometry as function of t

remark: certain quantities such as entanglement entropy are sensitive to **behind horizon physics** (away from equilibrium)

Translating to the CFT



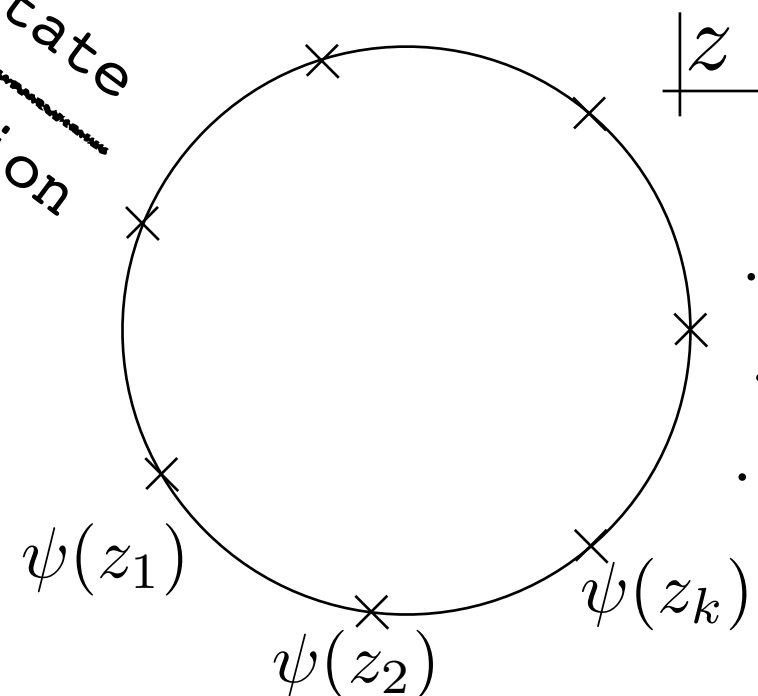
start in excited state at $t=0$:

prepare by Euclidean path integral
 \rightarrow regulator σ

$t = 0$

prepare state
 for t -evolution

$$|\mathcal{V}\rangle = \frac{1}{\mathcal{N}} \prod_{k=1}^n \psi(e_k, \bar{e}_k) |0\rangle$$



Interrogating the CFT

Start probing the physics via $2n + p$ correlations

$$G(1, 2, \dots, p) = \langle \mathcal{V} | \mathcal{Q}_1, \dots, \mathcal{Q}_p | \mathcal{V} \rangle$$

we want to approach smooth, semi-classical gravity

$$c \rightarrow \infty$$

$$n \rightarrow \infty$$

$$\sigma \rightarrow 0$$

$$E \sim nh_\psi/\sigma \rightarrow \mathcal{O}(c)$$

infinite-point correlations in strongly-coupled CFT!

Benefits of 2D CFT

in the semi-classical limit (large c), get sum of exponentials

$$G(1, 2, \dots p) = \sum_{\text{blocks}} a_k e^{-\frac{c}{6} f_k^{(n)}(1, 2, \dots p)}$$

correlator approximated by largest term, the identity block

“it from **id**”

*the dominant contribution comes from the identity Virasoro block, that is the unit operator **id** and all its descendants $T, \partial T, T^2, T\partial T, \dots$, (multi-graviton exchange in bulk)*

subleading corrections exponentially suppressed in $e^{-c} \sim e^{-1/G}$

still need to calculate the semi-classical block:

**CONFORMAL SCALAR FIELD ON THE HYPERELLIPTIC CURVE
AND CRITICAL ASHKIN-TELLER MULTIPOINT
CORRELATION FUNCTIONS**

AI.B. ZAMOLODCHIKOV

Scientific Council of "Cybernetics", Academy of Sciences, USSR

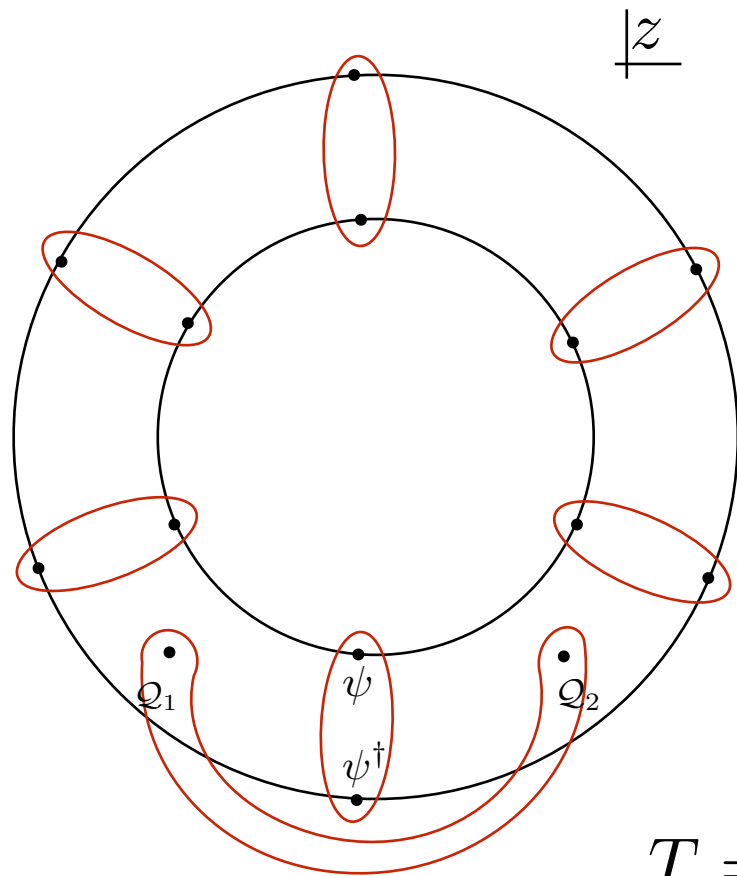
Received 3 December 1986

A multipoint conformal block of Ramond states of the two-dimensional free scalar field is calculated. This function is related to the free energy of the scalar field on the hyperelliptic Riemann surface under a particular choice of boundary conditions. Being compactified on the

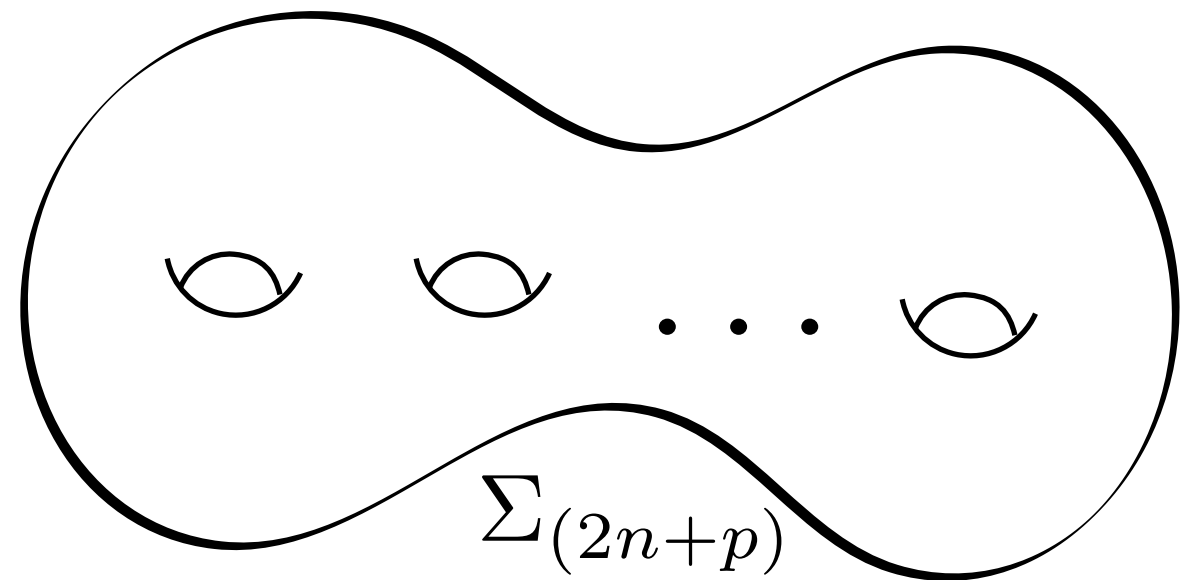
USSR: fighting hyper-intelligent Robot overlords with CFT?

The Monodromy Method

each contraction of operators in the plane defines a **cycle**



defines

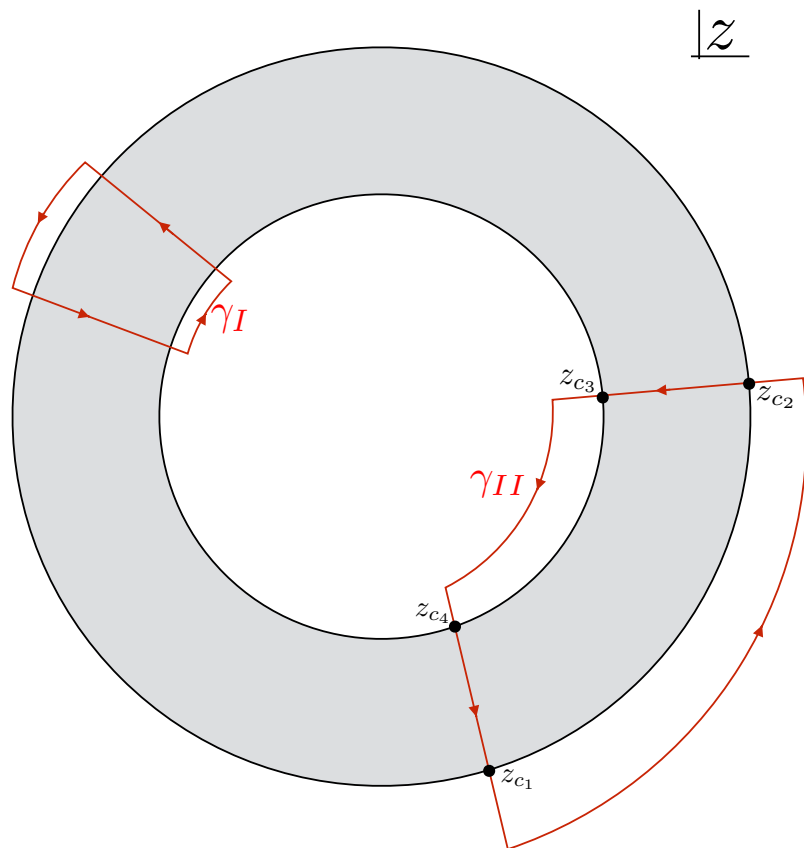


$$T = \sum_{k=1}^{2n+p} \left[\frac{6h/c}{(z - z_k)^2} + \frac{c_k}{z - z_k} \right]$$

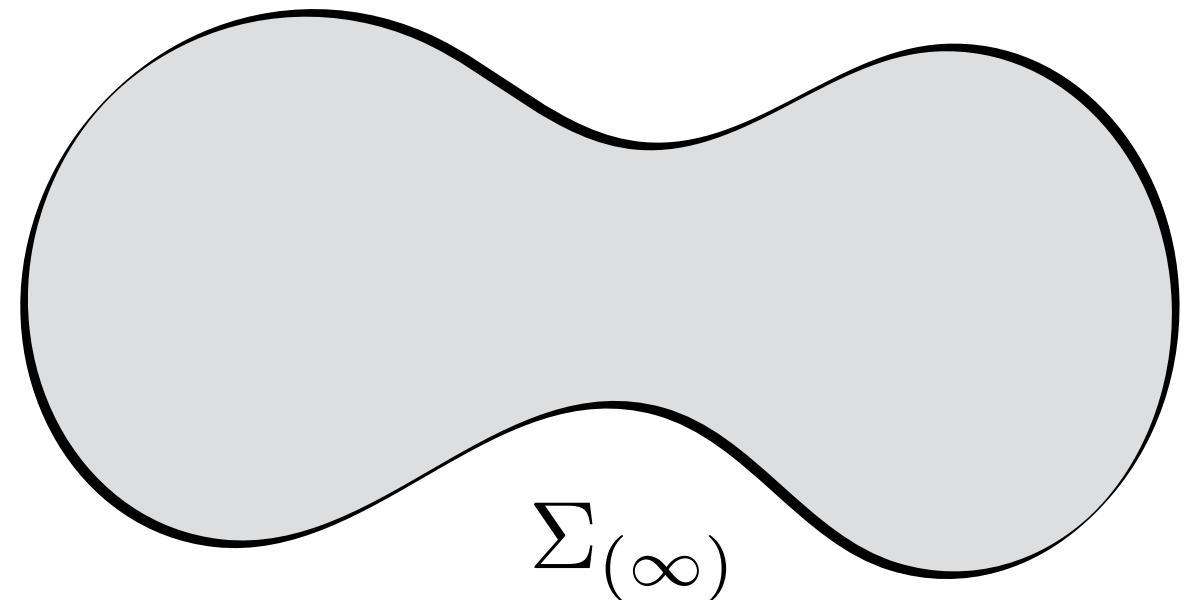
fix monodromies of $y''(z) + Ty(z) = 0 \xrightarrow{c_k} f_k^{(n)}(1, 2, \dots, p)$

Taking the smooth Limit

generally a hard problem, big simplification occurs for $n \rightarrow \infty$



defines



stress tensor \mapsto distribution on $\Sigma_{(\infty)}$

continuum monodromy method

$$f_0^\infty(1, 2, \dots, p)$$



3D semi-classical gravity

$$\mathcal{L}_{\text{geo}}(1, 2, \dots, p)$$

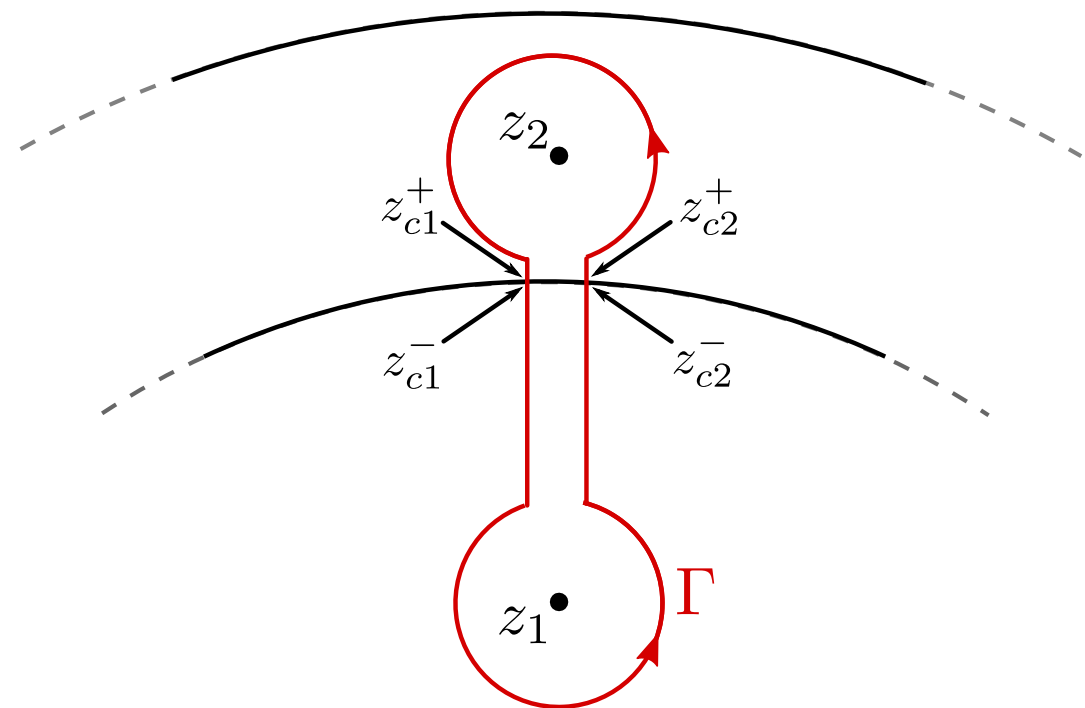
Two-point Autocorrelation

let us now return to the black hole and compute

$$G(t) = \text{tr} \rho \mathcal{O}(t) \mathcal{O}(0)$$

in the collapse state $|\mathcal{V}\rangle$

$$G(1, 2) =$$



can be done [analytically](#):

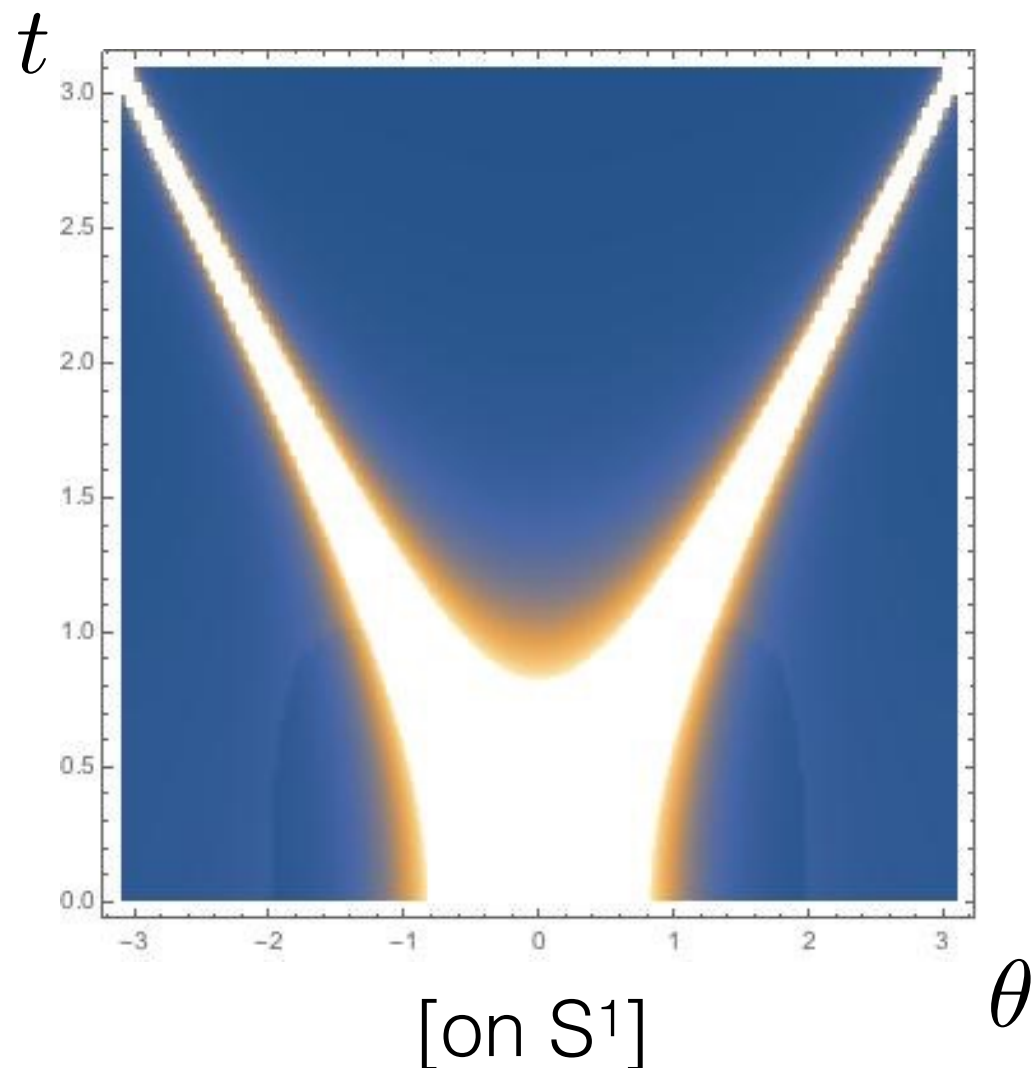
$$G(t_1, t_2) = \left(\frac{1}{\pi T} \cos \left(\frac{t_1}{2} \right) \sinh (\pi T t_2) - 2 \sin \left(\frac{t_1}{2} \right) \cosh (\pi T t_2) \right)^{-2\Delta_{\mathcal{Q}}}$$

General two-point function

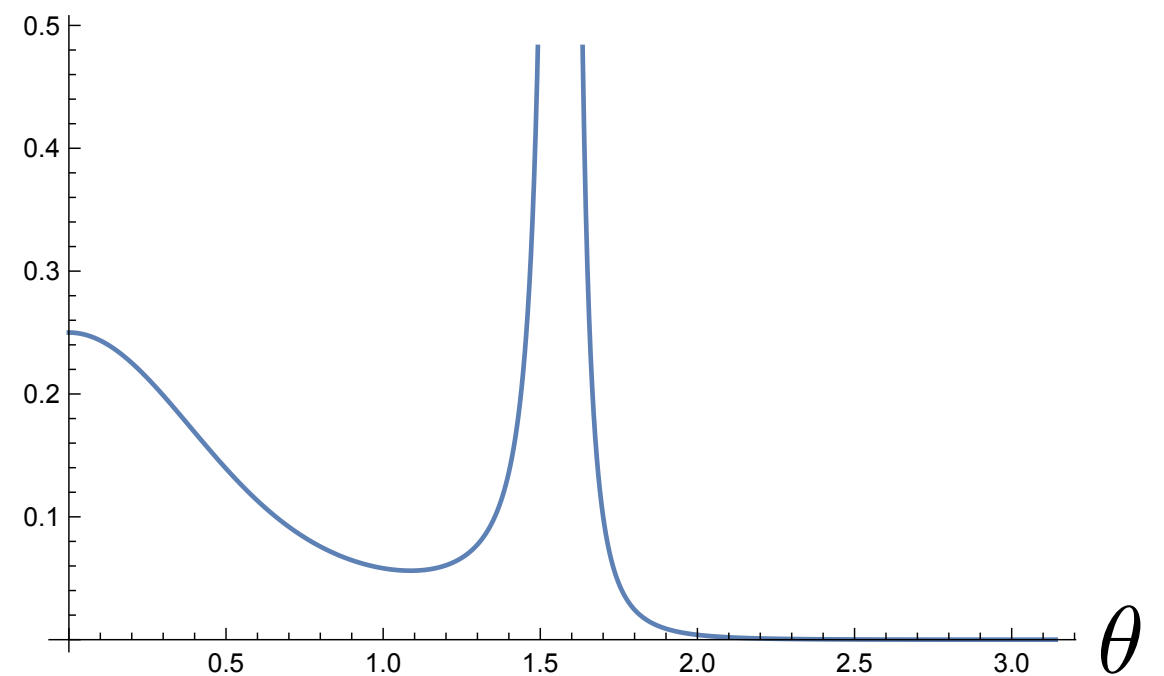
we are also interested in the general case

$$G(t, x) = \text{tr} \rho \mathcal{O}(t, x) \mathcal{O}(0)$$

some illustrative results:



$$\text{Re}G(\pi/2, \theta)$$



$$G(t, x) \sim e^{-2\pi T \Delta^{\mathcal{Q}} x} \quad [\text{on line}]$$

Physical Consequences

not (yet) known from gravity (but matches known limits)
 \implies CFT prediction for 3D gravity

The correlation function decays without bound at large time

$$G(t_1, t_2) \sim \exp\left(-\frac{2\pi\Delta^{\mathcal{Q}}t}{\beta}\right)$$

Manifestly in conflict with unitarity: **CFT loses information!**

Can also compute entanglement entropy of interval A

$$S(A) \rightarrow S_{\text{Gibbs}}(A; T) \longrightarrow \rho(A) = \rho_{\text{Gibbs}}(A; T)$$

Restoring Unitarity

This is the anti-information paradox: what happened to unitarity?

$$|G(t)| = \left| \sum_{n,k} e^{i(E_n - E_k)t} \Psi_n^*(\mathcal{V}) \langle n | \mathcal{Q} | k \rangle \langle k | \mathcal{Q} | \mathcal{V} \rangle \right| \neq 0$$

→ correlations **cannot become arbitrarily small** in $|\mathcal{V}\rangle$

Neglected contributions exponentially suppressed at $t=0$ (must be present due to crossing symmetry, e.g.)

$$\sum_{k \neq \text{vac}} a_k e^{-\frac{c}{6} f_k^\infty(1,2,\dots,p)} \sim e^{-S}$$

restore unitary at large time → non-perturbative effects in $1/G_N$

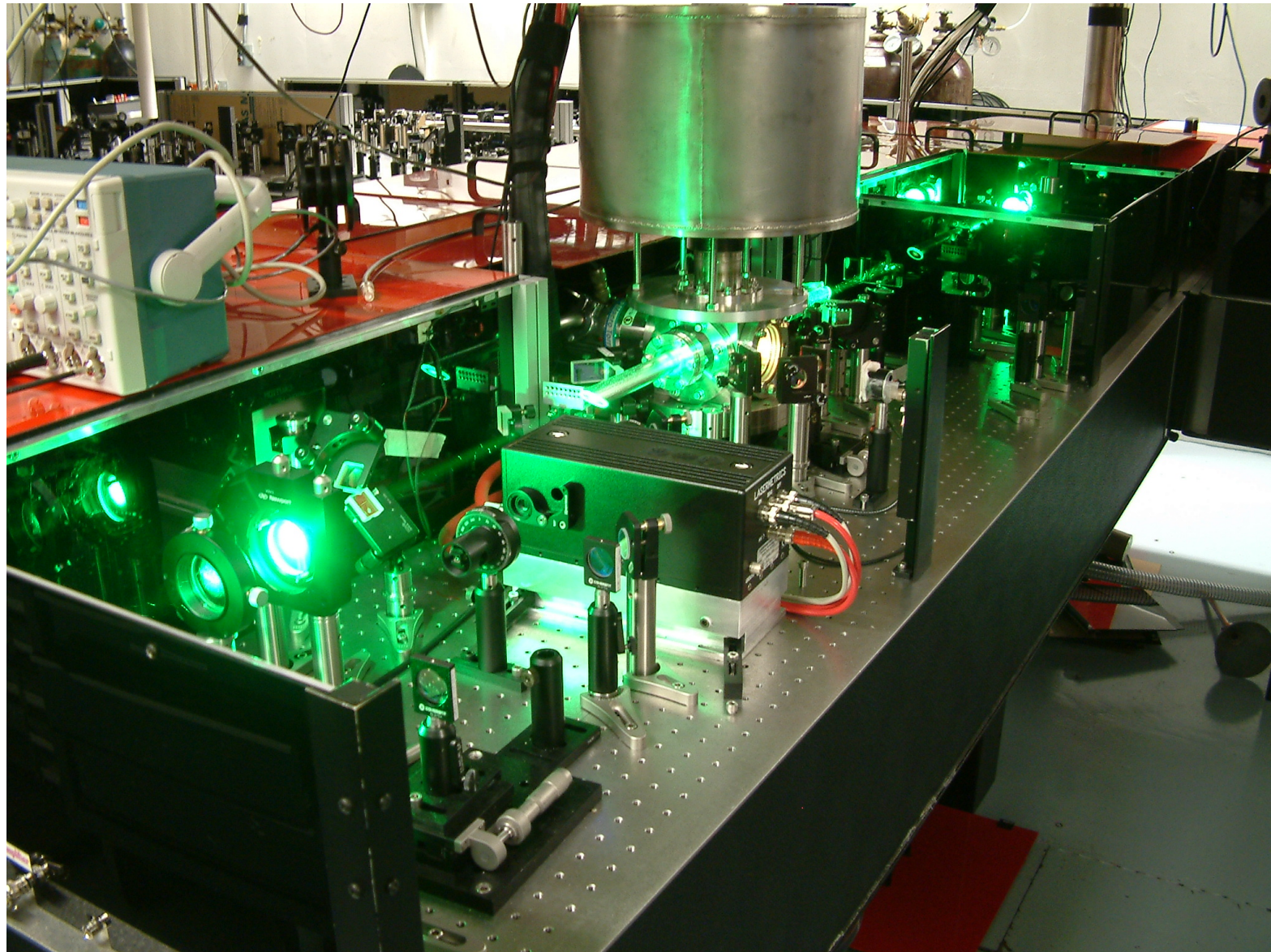
Conclusions

time-dependent 3D quantum gravity with matter in $1/c$ expansion
'it from id' \rightarrow ideal arena to think about quantum BHs

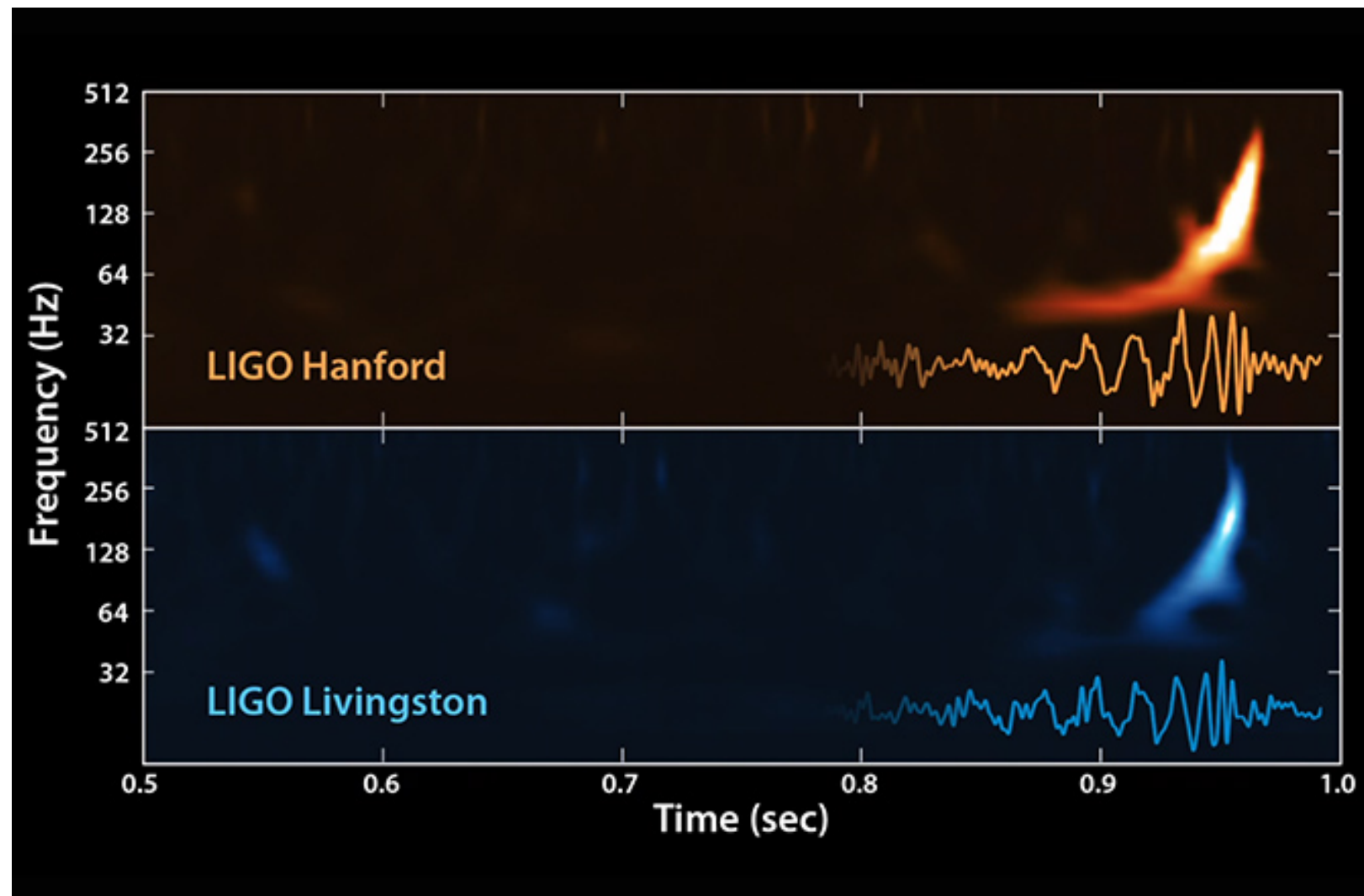
translates to detailed questions about thermalization in strongly-coupled CFT. New approach using monodromy (conformal blocks...)

correlation functions seemingly **violate unitarity** (naïve).
non-perturbative corrections in c restore unitarity

on gravity side these correspond to non-perturbative effects in G_N .
geometric interpretation? bulk interpretation?



AdS/CMT: quantum gravity in the lab



~~AdS/CMT: quantum gravity in the lab~~

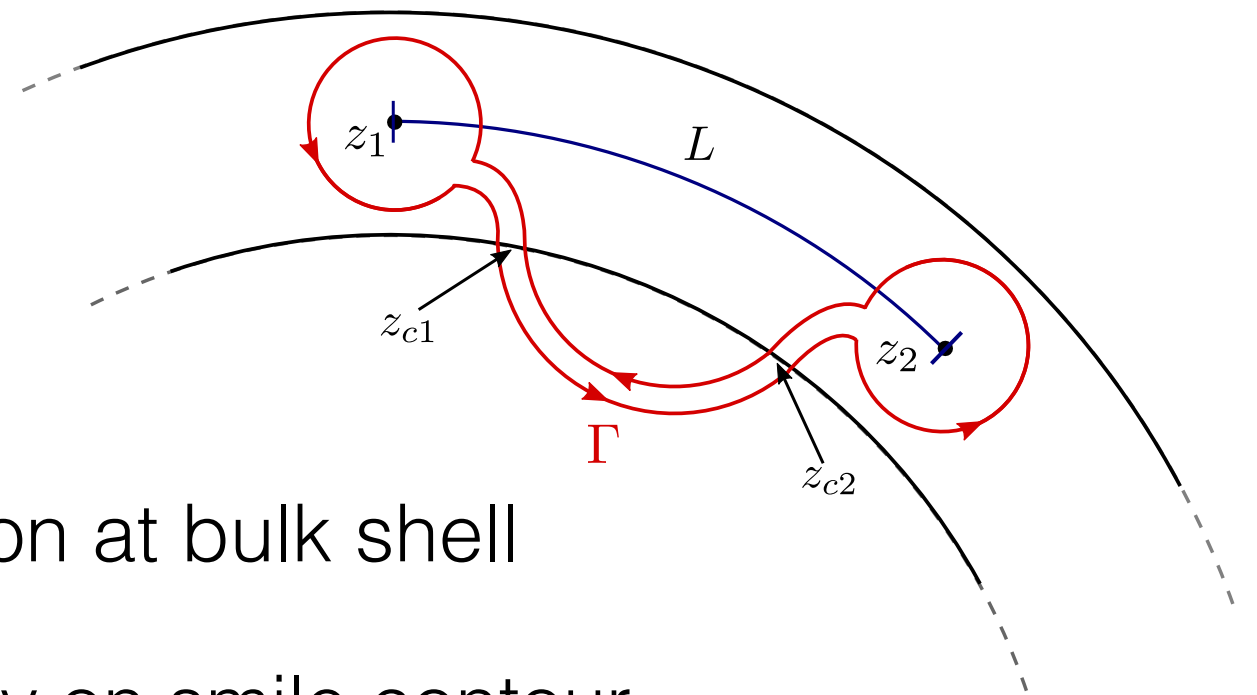
AdS/CMT: condensed matter in the Universe!

thank you!

entanglement entropy

Q-type operators \rightarrow twist insertions: $G_q(t) = \langle \mathcal{V} | \sigma_q(t, \ell_1) \tilde{\sigma}_q(t, \ell_2) | \mathcal{V} \rangle$

$$S(A) = \lim_{q \rightarrow 1} \frac{1}{1 - q} G_q(t)$$



crossing points z_{c1} & $z_{c2} \leftrightarrow$ refraction at bulk shell

it from id \rightarrow require trivial monodromy on smile contour

write $z_1 = e^{i\theta_1}$, $z_2 = e^{i(\theta_1 + L)}$ & continue to Lorentzian time $\theta_1 = t$

maximize $S(A)$ over crossing points \rightarrow parametric equation for $S(t)$

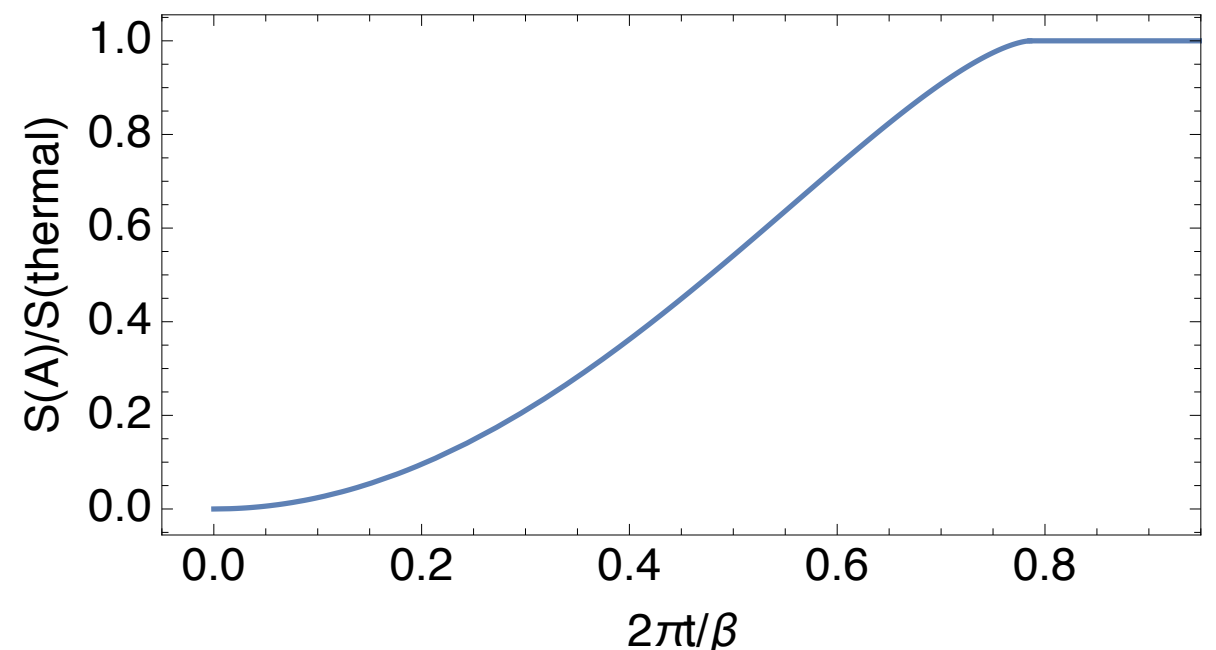
entanglement entropy

Implicit formula for growth of entanglement entropy:

$$t = \frac{\beta}{2\pi} \cosh^{-1} \left\{ \cosh(2\pi T q) + 2\pi T \tan\left(\frac{L}{2} - q\right) \sinh(2\pi T q) \right\}$$
$$S_{EE} = \frac{c}{3} \log \left\{ \frac{\sin\left(\frac{L}{2} - q\right) \cosh(2\pi T q) + \frac{1}{2\pi T} \left[1 + \frac{1}{2} \{ 1 + 4\pi^2 T^2 \} \tan^2\left(\frac{L}{2} - q\right) \right] \cos\left(\frac{L}{2} - q\right) \sinh(2\pi T q)}{\epsilon_{UV}/2} \right\}$$

matches **exactly** global AdS₃ Vaidya:

- thermal at late time
- EE growth = change of channel
- sees **beyond horizon**



CFT calculation shows that purity of state is preserved: $S(A) = S(A^c)$

alternative picture: IN-IN computation

2.) evolve in Lorentzian time until Q-operator insertion point(s)

1.) prepare initial state by Euclidean evolution for time σ

3,4.) evolve back in Lorentzian time, then Euclidean time to form conjugate

